

A colorful mirror solution to the strong CP problem

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2303.06156 [hep-ph]

The strong CP problem

$$\mathcal{L}_{\text{SM}} \supset \frac{g_s^2 \bar{\theta}}{32\pi^2} \epsilon^{\mu\nu\rho\sigma} G_{\mu\nu}^a G_{\rho\sigma}^a$$

Total derivative : invisible in perturbation theory

But non-perturbative effects !

In particular, **neutron electric dipole moment** (EDM)

[**Baluni '79, Crewther/Di Vecchia/Veneziano/Witten '79**]

$$\mathcal{L}_{\text{EDM}} \supset \frac{id_n}{2} \bar{n} \gamma_5 \gamma_{\mu\nu} n F^{\mu\nu}$$

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 $\approx 10^{-2} \bar{\theta} e \text{ GeV}^{-1}$

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measured to be
 $\lesssim 10^{-12} e \text{ GeV}^{-1}$

[Pendlebury et al '15]

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But non-perturbative

In particular, **neutral**

Strong CP problem :

$$\bar{\theta} \lesssim 10^{-10}$$

[**Banerjee '79, Crewther, Di Vecchia, Veneziano/Witten '79**]

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The strong CP problem

$$\mathcal{L}_{\text{SM}} \supset \frac{g_s^2 \bar{\theta}}{32\pi^2} \epsilon^{\mu\nu\rho\sigma} G_{\mu\nu}^a G_{\rho\sigma}^a$$

Even worse : valid in the basis where **all fermion masses are real**

Before EW symmetry breaking :

$$\bar{\theta} = \theta + \arg \det(Y_u Y_d)$$

$$\begin{aligned} \mathcal{L}_{\text{SM}} &\supset \frac{g_s^2 \theta}{32\pi^2} \epsilon^{\mu\nu\rho\sigma} G_{\mu\nu}^a G_{\rho\sigma}^a \\ &+ \bar{Q}_L Y_u u_R \tilde{H} + \dots \end{aligned}$$

Two unrelated quantities cancel almost perfectly !

Solutions to the strong CP problem

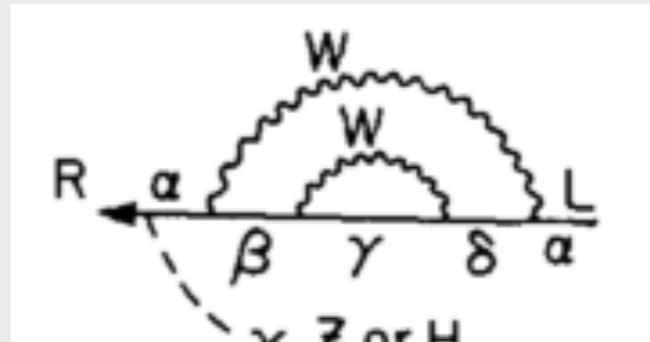
$$\mathcal{L}_{\text{SM}} \supset \frac{g_s^2 \bar{\theta}}{32\pi^2} \epsilon^{\mu\nu\rho\sigma} G_{\mu\nu}^a G_{\rho\sigma}^a$$

How to get $\bar{\theta} \approx 0$?

An effective field theorist's nightmare !

- non-decoupling contributions at all scales
- barely regenerated by renormalization group flow

[Ellis/Gaillard '79]



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How to get $\bar{\theta} \approx 0$?

An effective field theorist's nightmare ! IR solutions.

Make it unphysical

Make it dynamical

$$\bar{\theta} = \theta + \arg \det(Y_u Y_d)$$

Massless quark :
ambiguous !

['t Hooft '76]

Chiral sym. breaking : spurion $\theta - \langle \eta' \rangle$

Relaxes to zero !

[di Vecchia/Veneziano '80, Vafa/Witten '84]

Disfavored by lattice...

[Aoki et al '16]

$$\mathcal{L} \supset \frac{g_s^2}{16\pi^2} \frac{a}{f_a} G \tilde{G}$$

[Peccei/Quinn '77,
Weinberg '78, Wilczek '78]

Spurion $\bar{\theta} - \langle a \rangle / f_a$

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**Non-minimal
model building**

[Peccei/Quinn '77,
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Make it zero

By symmetry: **(C)P** !

Parity solutions to the strong CP problem

(C)P is not a symmetry of the SM ! Spontaneous breaking

CP : $\langle \phi \rangle \in \mathbb{C}$

[Nelson '84, Barr '84]

Parity solutions to the strong CP problem

P is not a symmetry of the SM ! Spontaneous breaking

$$Q_L(\mathbf{3}, \mathbf{2}, 1/6)$$

$$u_R(\mathbf{3}, \mathbf{1}, 2/3)$$

$$d_R(\mathbf{3}, \mathbf{1}, -1/3)$$

Parity solutions to the strong CP problem

P is not a symmetry of the SM ! Spontaneous breaking

New fermions ? New gauge fields ?

$$Q_L(3, 2, 1/6)$$

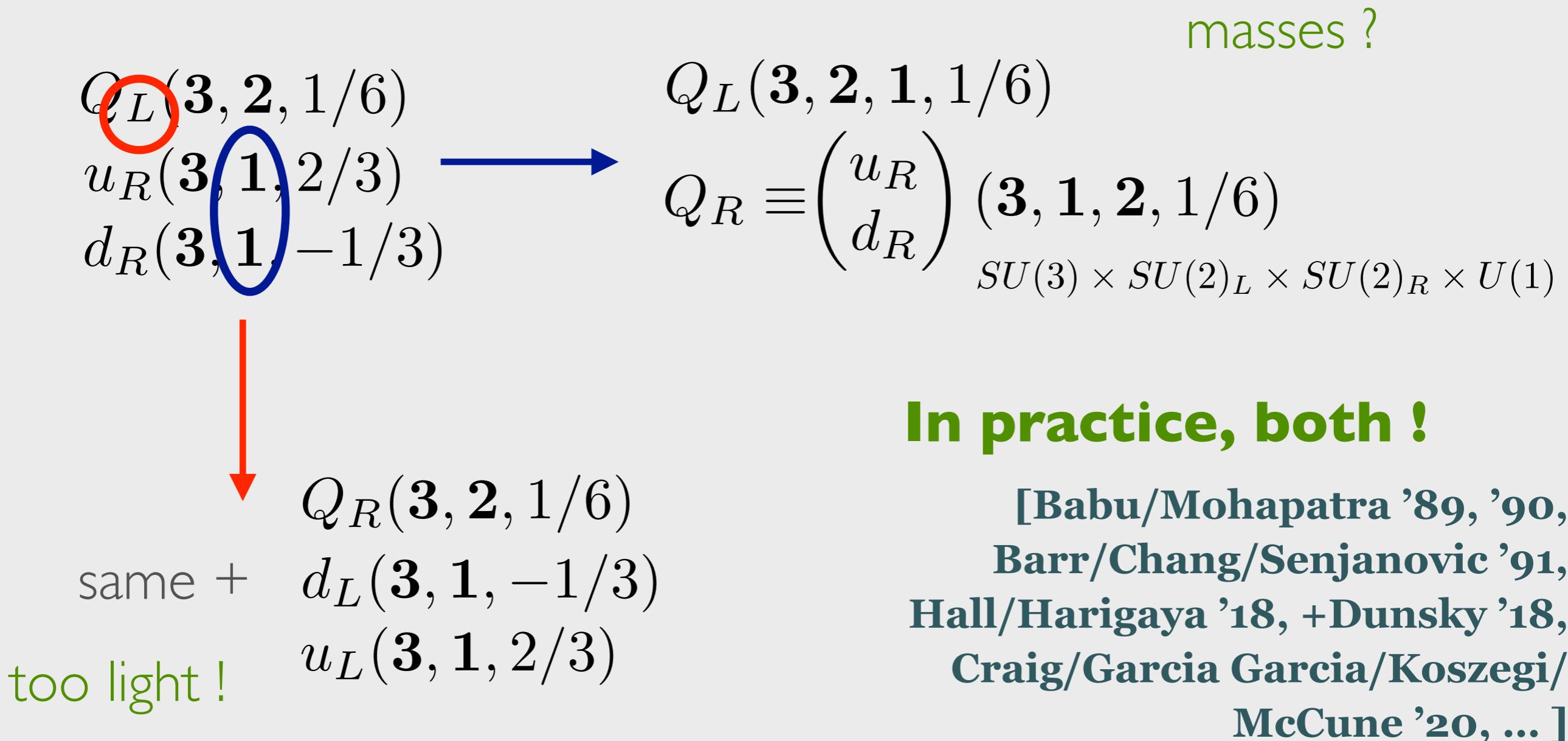
$$u_R(3, 1, 2/3)$$

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Parity solutions to the strong CP problem

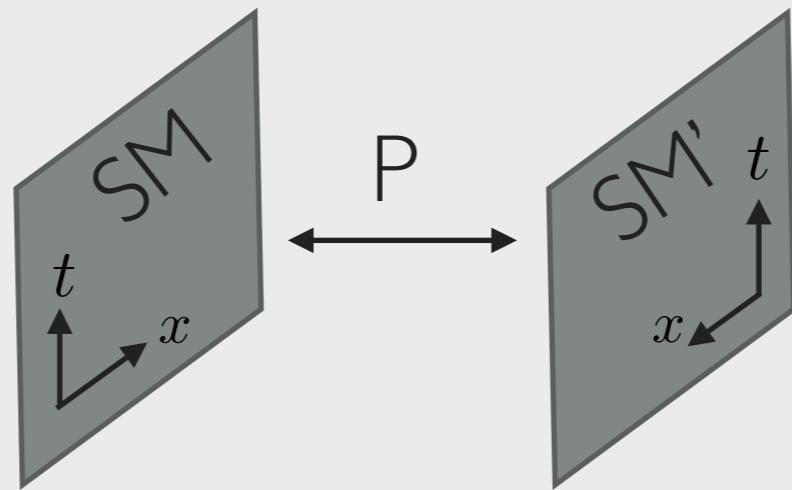
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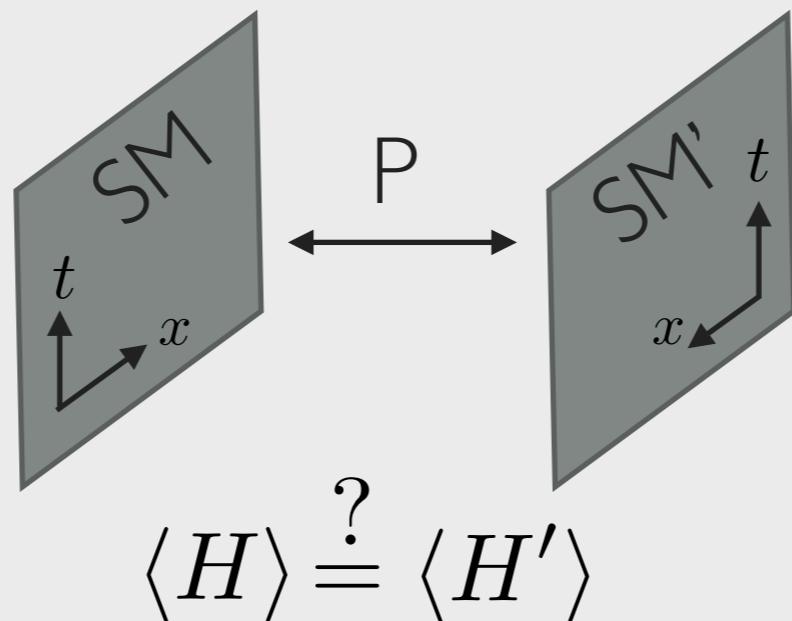
Parity solutions to the strong CP problem

Mirror world



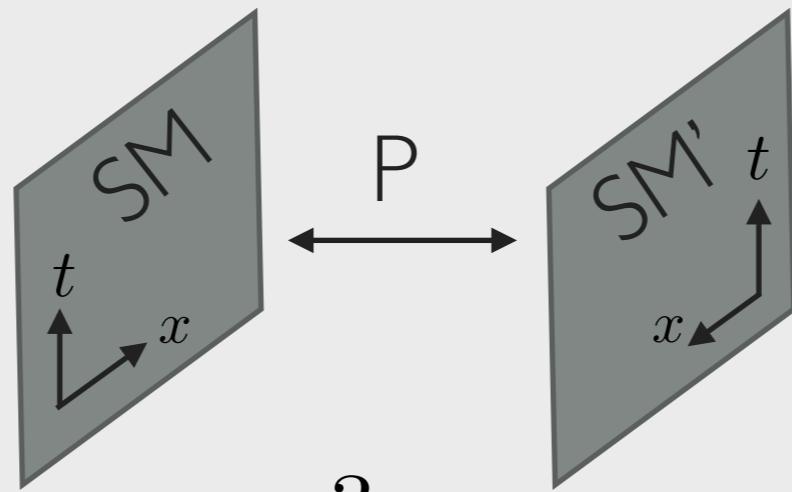
Parity solutions to the strong CP problem

Mirror world



Parity solutions to the strong CP problem

Mirror world and strong CP

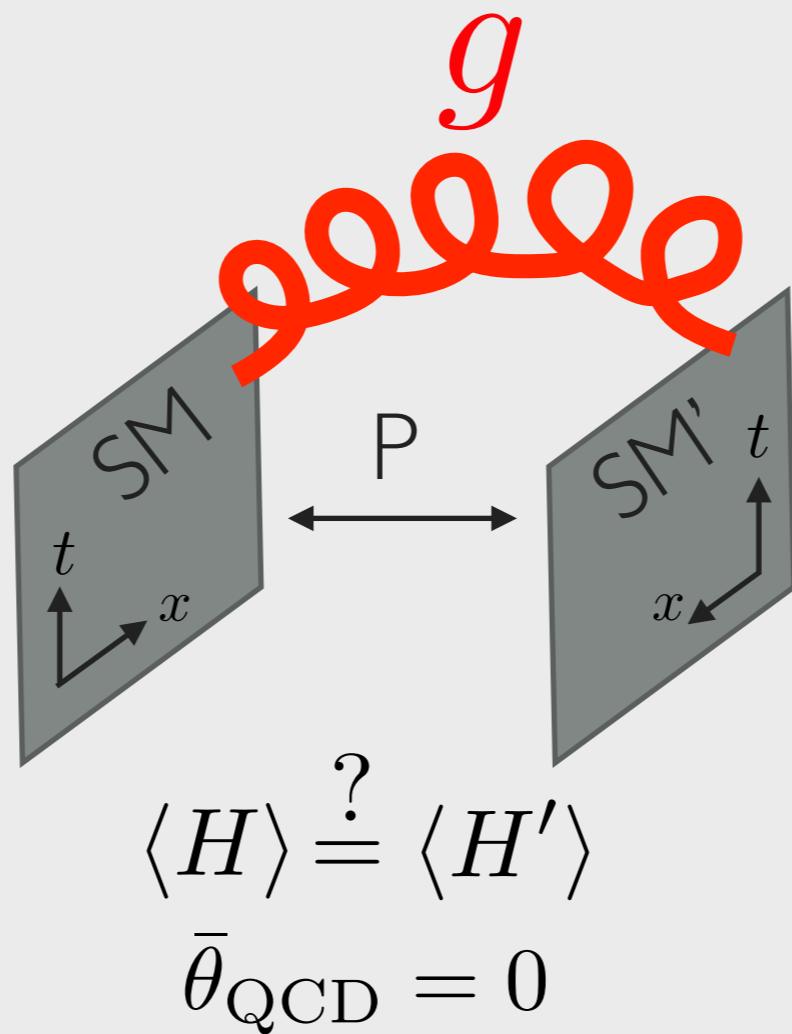


$$\langle H \rangle \stackrel{?}{=} \langle H' \rangle$$

$$\bar{\theta}' = -\bar{\theta}$$

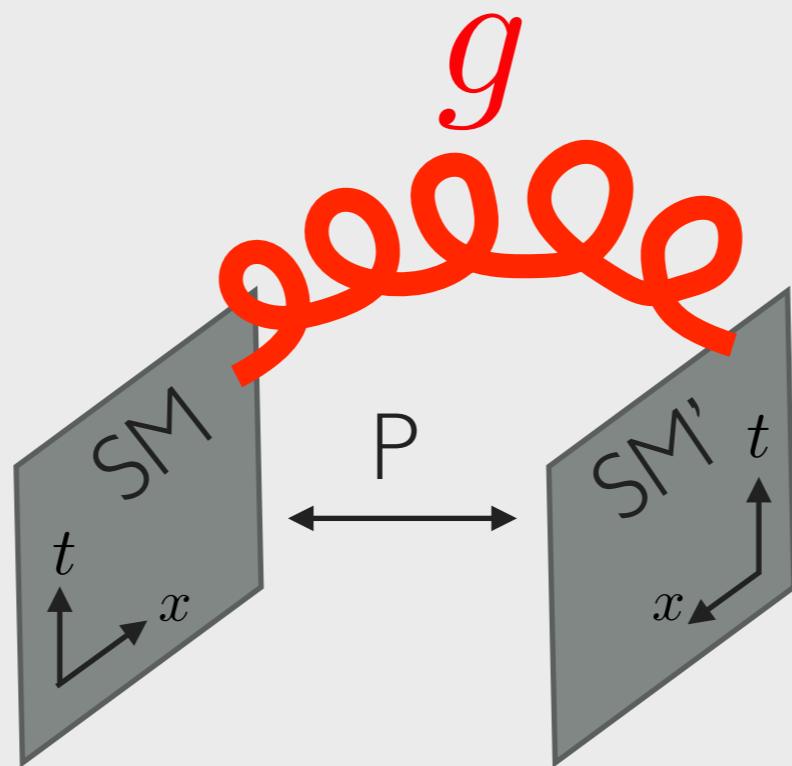
Parity solutions to the strong CP problem

Mirror world and strong CP. Need shared color (P-invariant on its own) !



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Mirror world and strong CP. Need shared color (P-invariant on its own) !

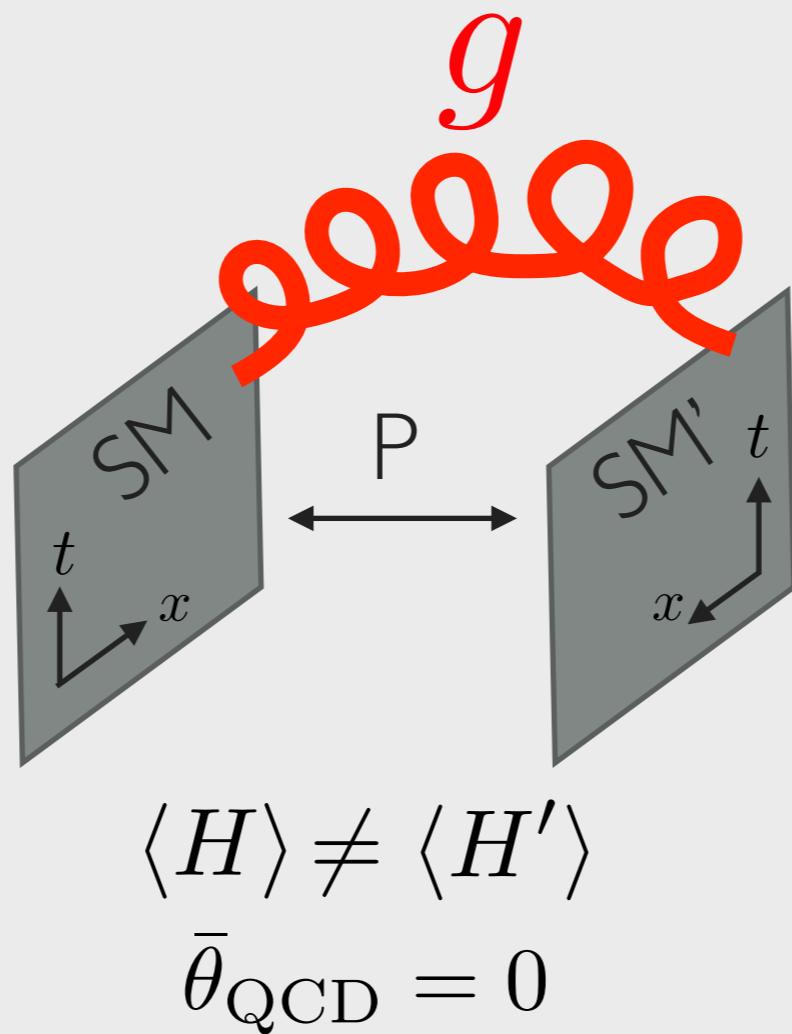


$$\langle H \rangle \neq \langle H' \rangle$$

$$\bar{\theta}_{\text{QCD}} = 0$$

Parity solutions to the strong CP problem

Mirror world and strong CP. Need shared color (P-invariant on its own) !



Few parts of that landscape are explored !

Parity solutions to the strong CP problem

- ? Need $\bar{\theta} \approx 0$ even **below the scale of parity breaking**
 - Similar to the SM case [Ellis/Gaillard '79]
 - Not in all extensions [de Vries/Draper/Patel '21]
 - Easier in mirror models [Barr/Chang/Senjanovic '91]

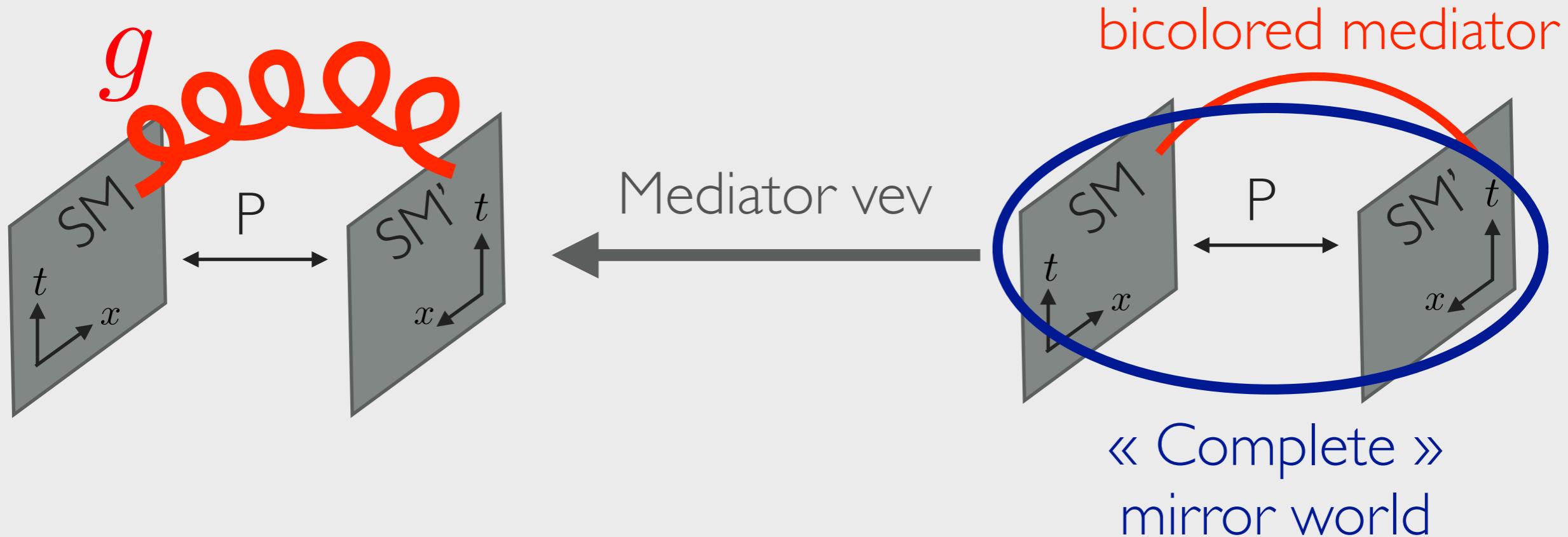
Our parity solution to the strong CP problem

We notice that



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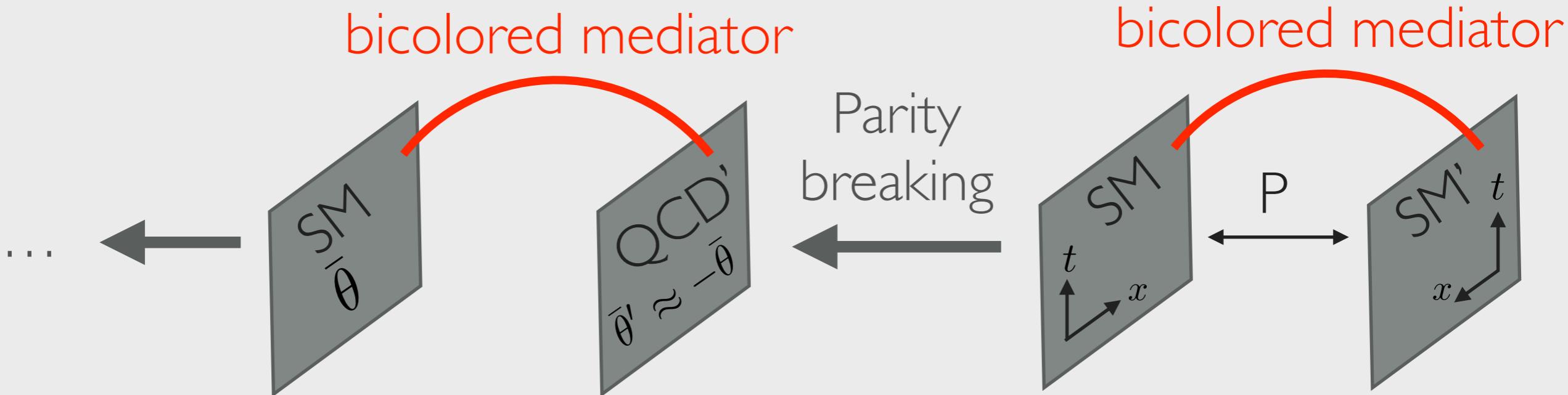


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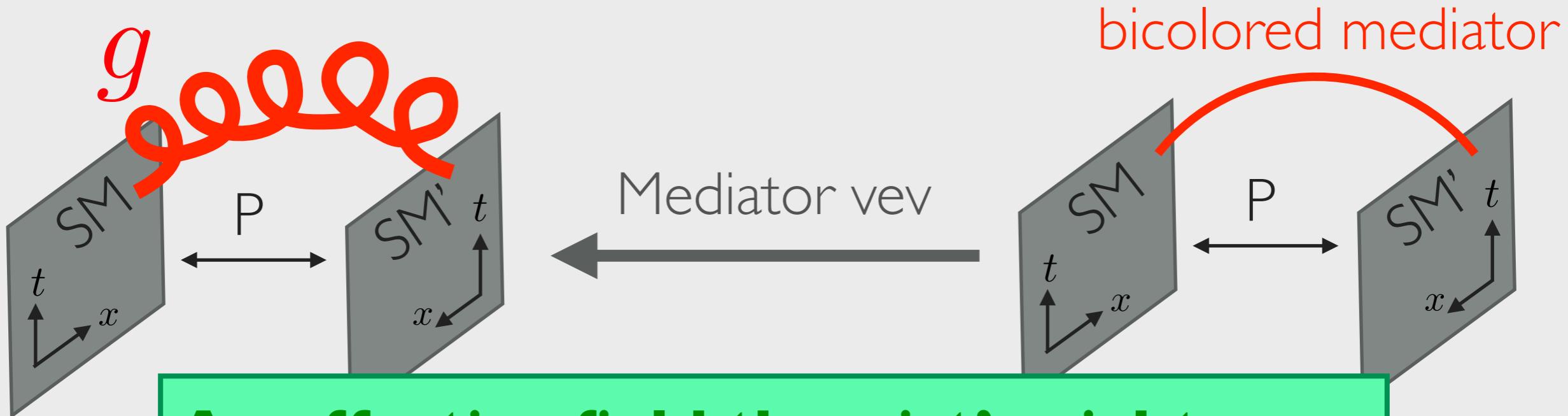


With the same starting point :

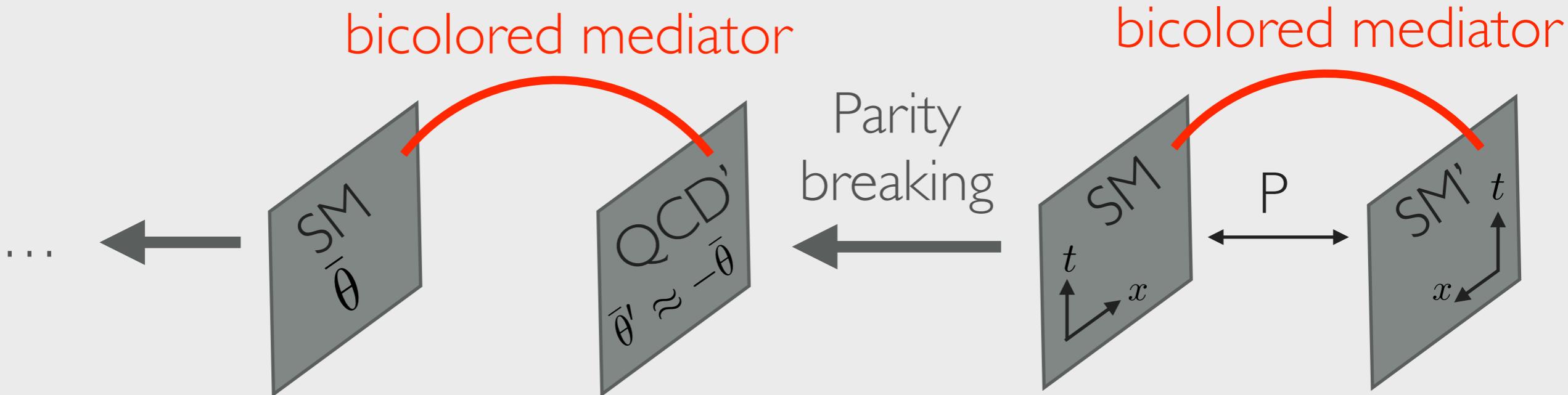


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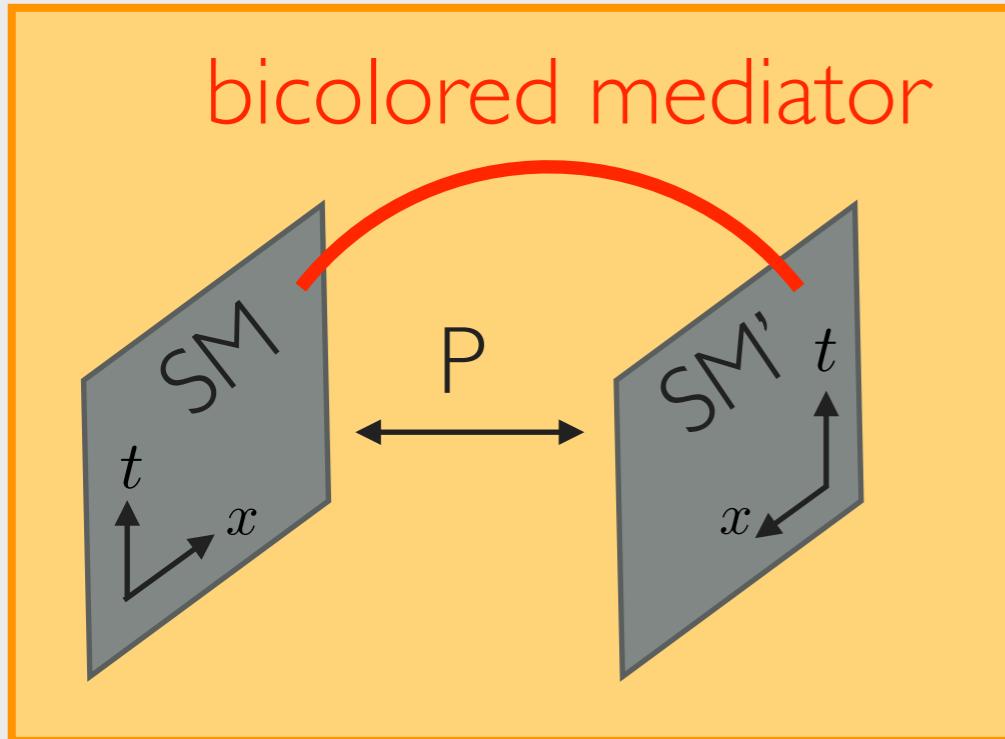
We notice that



With the same starting point :



Our parity solution to the strong CP problem



Bicolored mediator here:
bifundamental order
parameter $\langle \Sigma \rangle$

$$\langle \Sigma \rangle \propto v_3 \mathbf{1}$$

$$\implies g_3^2 G \tilde{G} = \Big|_{\text{along QCD}} g_3'^2 G' \tilde{G}'$$

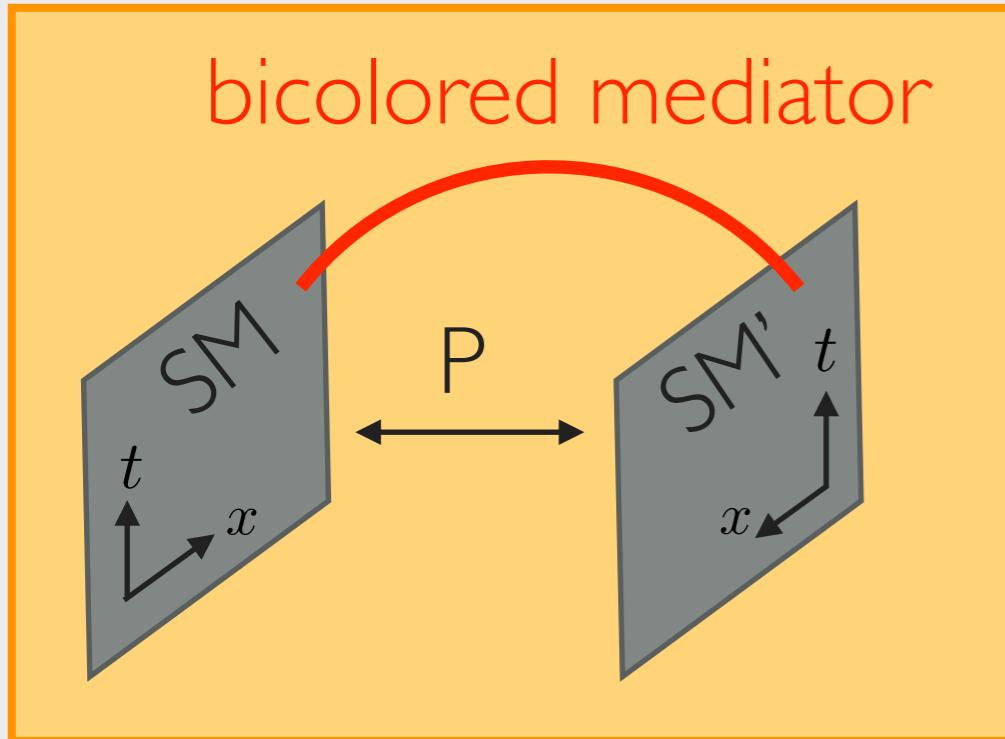
	$SU(3)$	$SU(2)_L$	$U(1)_Y$	$SU(3)'$	$SU(2)'$	$U(1)'$
Q	3	2	1/6	1	1	0
u^c	$\bar{3}$	1	-2/3	1	1	0
d^c	$\bar{3}$	1	1/3	1	1	0
L	1	2	-1/2	1	1	0
e^c	1	1	-1	1	1	0
H	1	2	1/2	1	1	0
Q'	1	1	0	$\bar{3}$	2	-1/6
u'^c	1	1	0	3	1	2/3
d'^c	1	1	0	3	1	-1/3
L'	1	1	0	1	2	1/2
e'^c	1	1	0	1	1	1
H'	1	1	0	1	2	-1/2



$$\left. \begin{aligned} \theta' &= -\theta \\ Y_q &= Y_{q'}^\dagger \end{aligned} \right\} \implies \bar{\theta}' = -\bar{\theta}$$

$$\implies \bar{\theta}_{\text{QCD}} = 0$$

Our parity solution to the strong CP problem



Below v_3 , colored mirror
quarks : need $\langle H' \rangle \gg \langle H \rangle$
(hence P-breaking) $\equiv_{v'}$

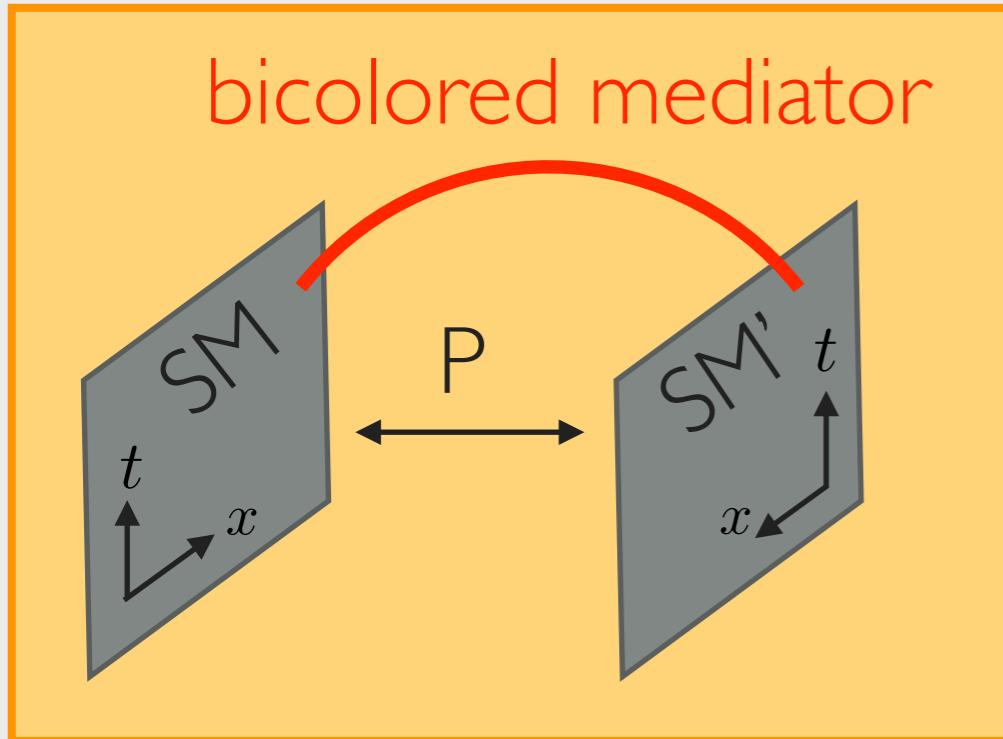
i.e. $v' \gtrsim 10^9$ GeV

	$SU(3)$	$SU(2)_L$	$U(1)_Y$	$SU(3)'$	$SU(2)'$	$U(1)'$
Q	3	2	$1/6$	1	1	0
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H	1	2	$1/2$	1	1	0
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$$Y_q = Y_{q'}^\dagger \implies \frac{m_q}{m_{q'}} = \frac{\langle H \rangle}{\langle H' \rangle}$$

Our parity solution to the strong CP problem



Below v_3 , colored mirror quarks : need $\langle H' \rangle \gg \langle H \rangle$
 (hence P-breaking) $\equiv_{v'}$

Achieved through soft breaking or radiative corrections

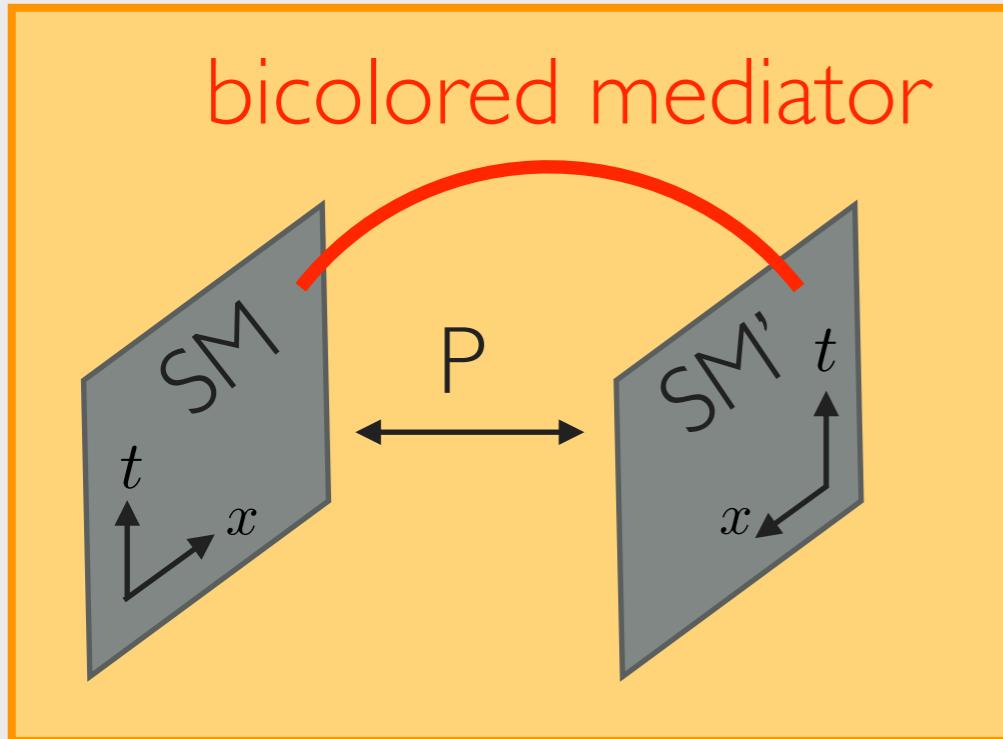
[Babu/Mohapatra '89,
 Hall/Harigaya '18]

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Very predictive model, two BSM scales : v_3 and v'

Different pheno on the parameter space. For $v_3 \ll v'$, **colored bosons** as lightest BSM states !

$$Y_q = Y_{q'}^\dagger \implies \frac{m_q}{m_{q'}} = \frac{\langle H \rangle}{\langle H' \rangle}$$

Our parity solution to the strong CP problem

Bicolored mediator : a **scalar** or **strongly interacting fermions**.

- Σ in $(\mathbf{3}, \mathbf{3}')$ of $SU(3) \times SU(3)'$ with potential

$$\begin{aligned} V(\Sigma) = & -m^2 \text{Tr}(\Sigma \Sigma^\dagger) + c \text{Tr}^2(\Sigma \Sigma^\dagger) \\ & + \tilde{c} \text{Tr}(\Sigma \Sigma^\dagger)^2 + (\tilde{m} \det(\Sigma) + h.c.) \end{aligned}$$

$\mathbf{3}$ or $\bar{\mathbf{3}}$

Breaking to the diagonal $SU(3)$ in a large fraction of parameter space (but no (C)P breaking)

[Bai/Dobrescu '17]

-

	$SU(N)$	$SU(N)'$	$SU(3)$	$SU(3)'$	Breaking to the diagonal $SU(3)$ à la technicolor
ψ_L	N	1	3	1	
ψ_R	N	1	1	3'	
ψ'_L	1	N	$\bar{3}$	1	
ψ'_R	1	N	1	$\bar{3}'$	[Weinberg '76, Susskind '78]

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ψ_L	N	1	3	1
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ψ'_L	1	N	3	1 P for 3' = 3
ψ'_R	1	N	1	3' P for 3' = 3

Diagram showing the representation assignments for the fields under $SU(N) \times SU(N)' \times SU(3) \times SU(3)'$. The fields are:
 ψ_L : N, 1, 3, 1
 ψ_R : N, 1, 1, 3'
 ψ'_L : 1, N, 3, 1 P for $3' = 3$
 ψ'_R : 1, N, 1, 3' P for $3' = 3$

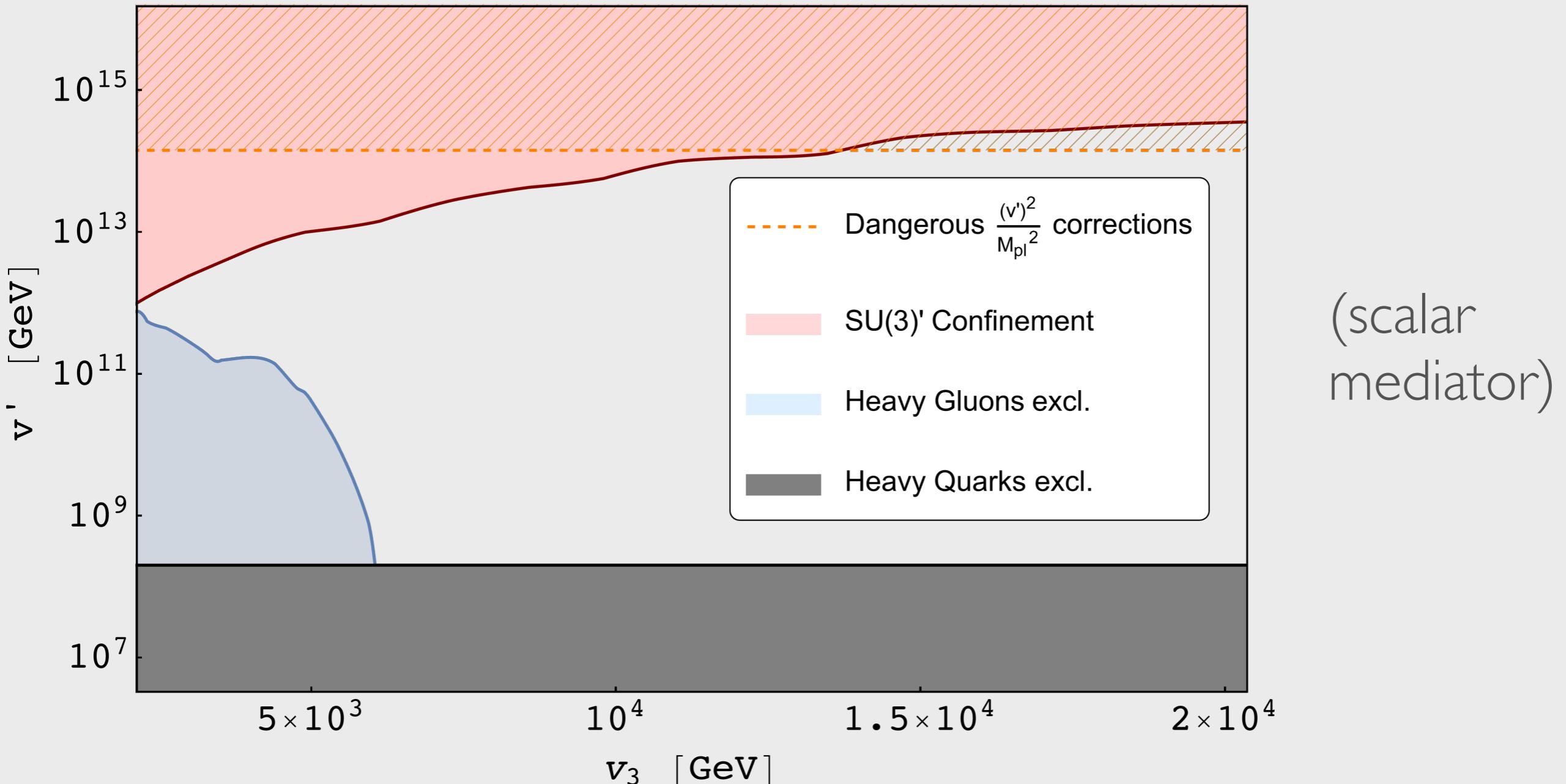
A circular arrow diagram indicates a symmetry between the $SU(3)'$ representations 1 and 3'. Labels "P for $3' = 3$ " are placed near the arrows.

Our parity solution to the strong CP problem

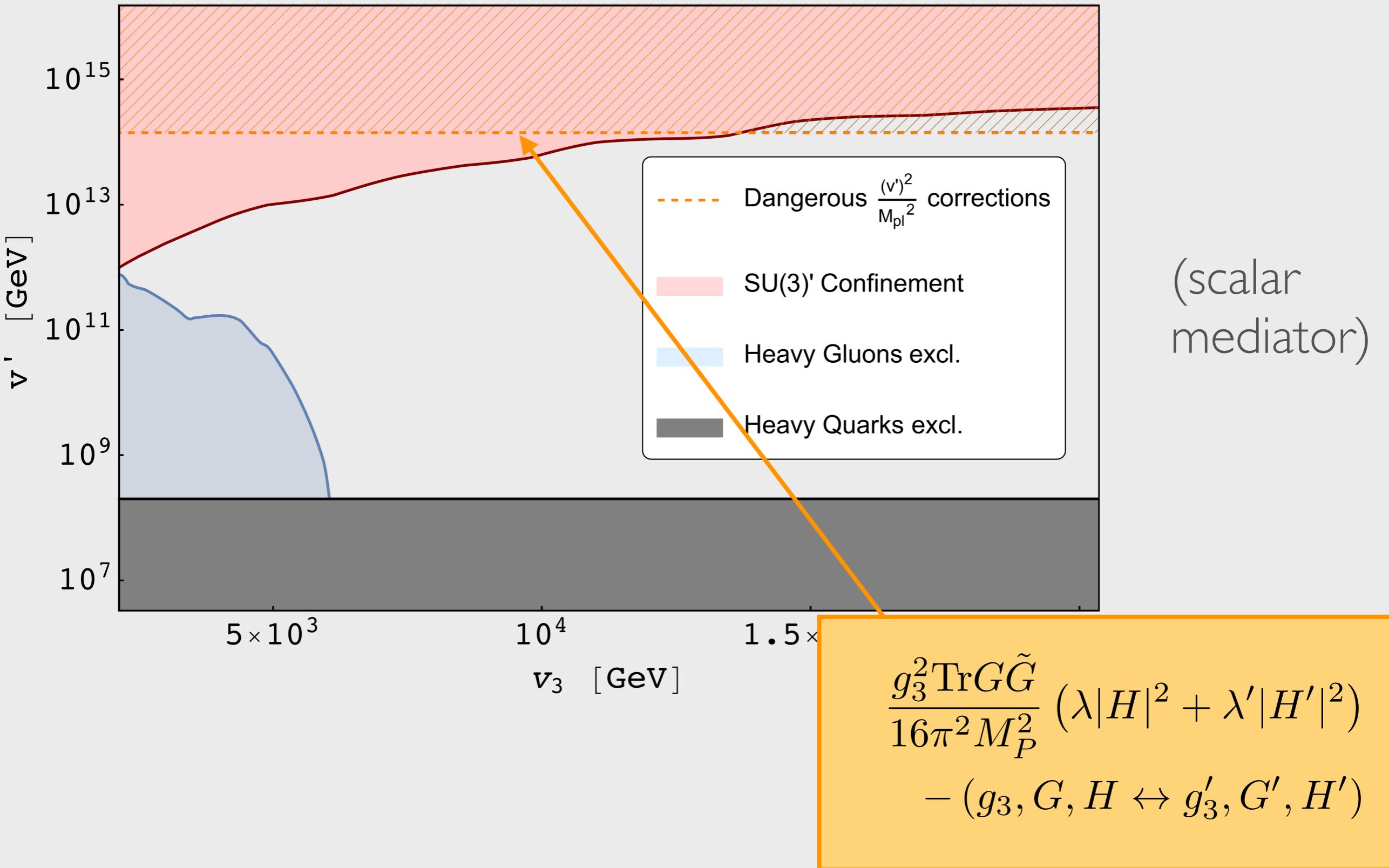
- ? Need $\bar{\theta} \approx 0$ even **below the scale of parity breaking**

Only mediators: gluons, bicolor mediator or heavy Higgs. Only CP phase: CKM. **Very small contributions** (at least 3-loops) **to** $\bar{\theta}_{\text{QCD}}$

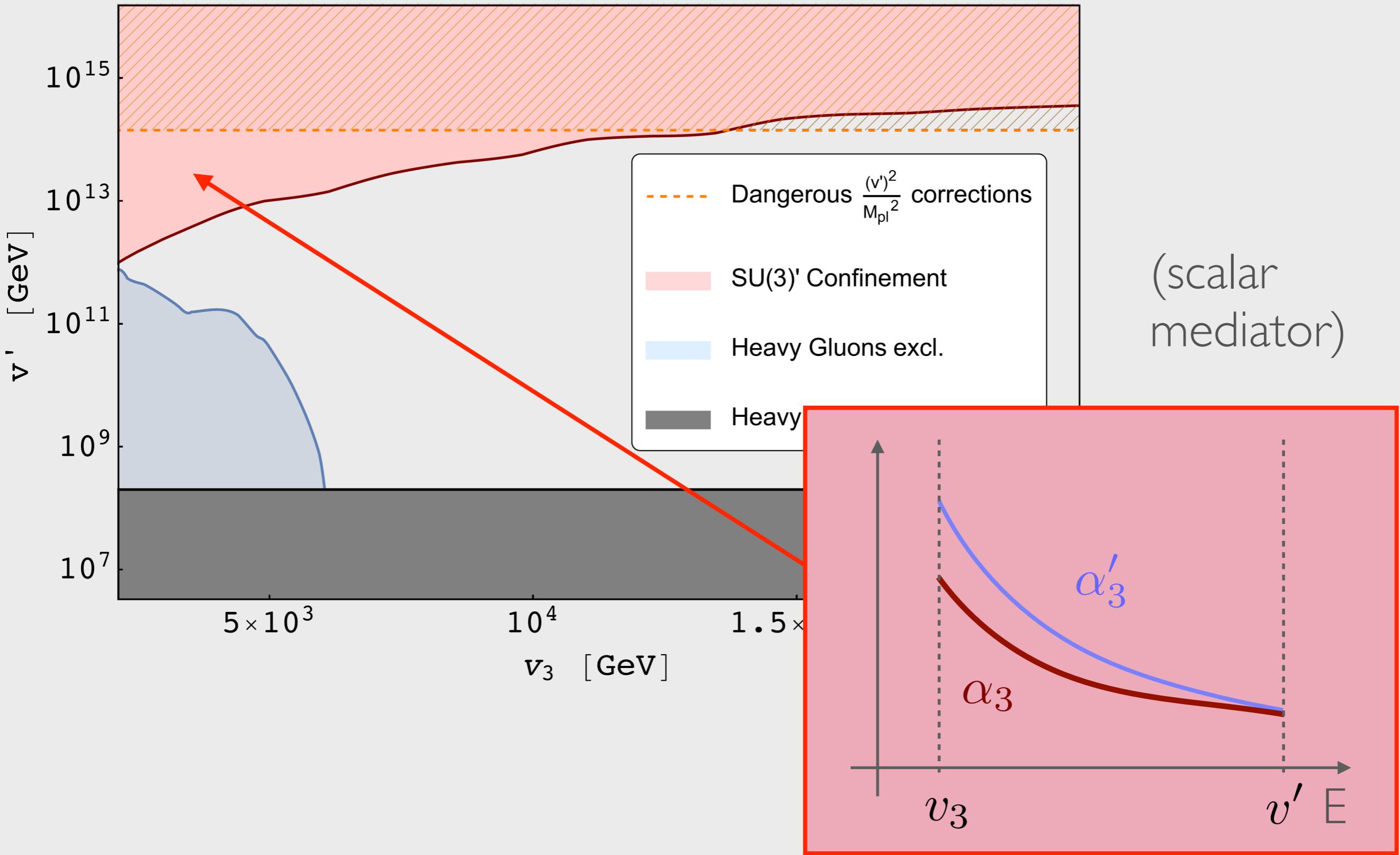
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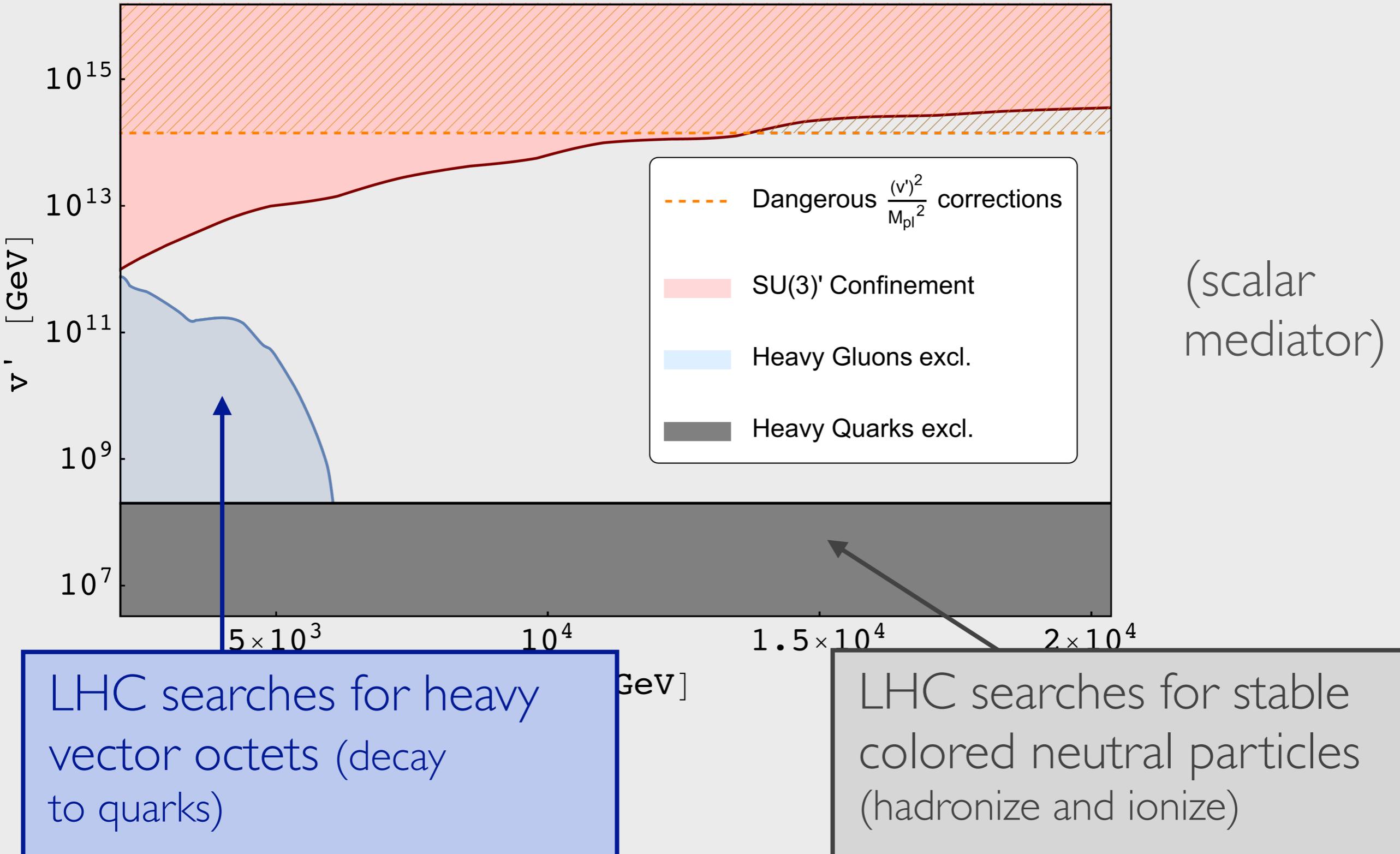
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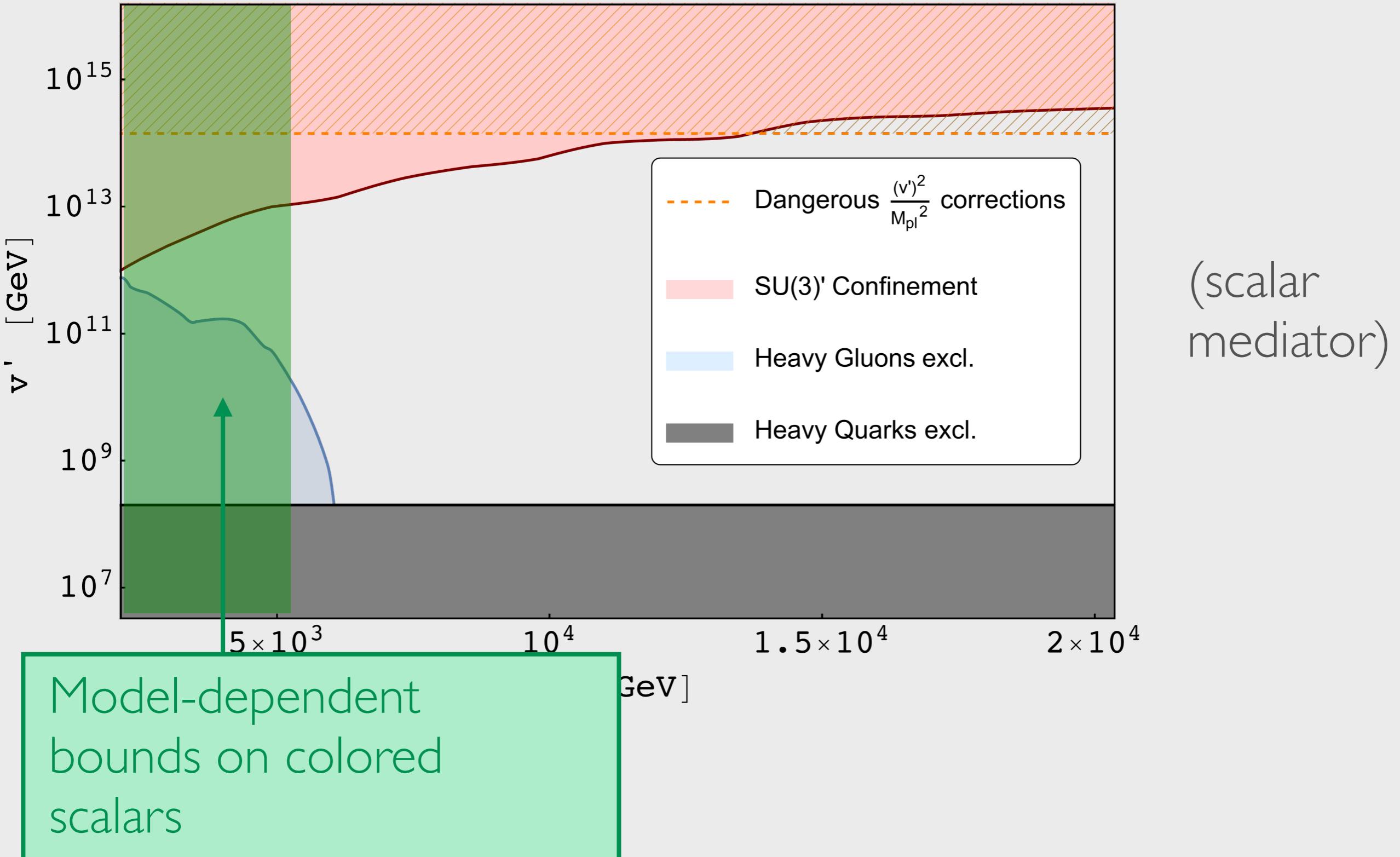
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Our parity solution to the strong CP problem



Outlook

First study of a **parity solution to the strong CP problem in a « complete » mirror world**

Corrections to $\bar{\theta}_{\text{QCD}}$ very much under control

Few free parameters, but quite different physics in the parameter space !

Colored bosons could be the first accessible signal

More to explore ! Landscape of models still quite uncharted. In this model: interesting (and new) cosmology, with possible dark matter candidates, phase transitions, ...

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THANK YOU