A one-parameter refinement of the Razumov-Stroganov correspondence





LPTM, Cergy-Pontoise

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Based on joint work with A. Sportiello (Milan University)

Outline

- 1 The Temperley-Lieb Stochastic process
- 2 Alternating Sign Matrices or Fully Packed Loop configurations
- 3 The Razumov Stroganov (ex-)conjecture
- One parameter refinement of the Razumov–Stroganov correspondence
- 5 Conclusions & Open Problems

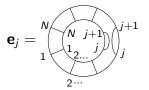
Recall the definition of the *Cyclic Temperley-Lieb Algebra* $\mathrm{CTL}_N(\tau)$: free algebra with generators $\{\mathbf{e}_i\}_{i\in\mathbb{Z}}$ and the rotation R

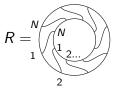
$$\mathbf{e}_{i} = \mathbf{e}_{i+N} \qquad \qquad R^{\pm 1} \mathbf{e}_{i} = \mathbf{e}_{i\pm 1} R^{\pm 1}$$

$$\mathbf{e}_{i} = \tau \mathbf{e}_{i} \qquad \qquad \mathbf{e}_{i} \mathbf{e}_{i\pm 1} \mathbf{e}_{i} = \mathbf{e}_{i}$$

$$[\mathbf{e}_{i}, \mathbf{e}_{j}] = 0 \qquad \qquad |i - j| \neq 1 \pmod{N}.$$

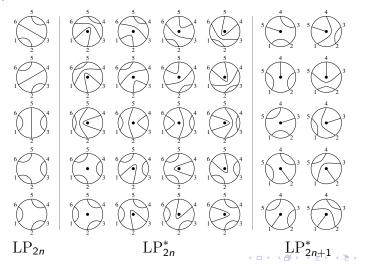
Graphical representation





Action on link patterns

We shall be interested in three kind of representations of $\mathrm{CTL}_N(\tau)$ on *link patterns*



For
$$au=1$$
 the operator

$$H_N = \frac{1}{N} \sum_{i=1}^N \mathbf{e}_i$$

is the markov Matrix of the so called Temperley-Lieb Stochastic process [Batchelor, de Gier & Nienhuis, Razumov, Stroganov].

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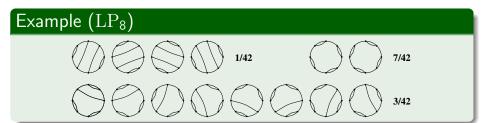
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ullet We are interested in its *Stationary Probability* for the representations LP_{2n} and LP_N^*

$$|\Psi_n\rangle := \sum_{\pi \in \operatorname{LP}_{2n}} \Psi_n(\pi) |\pi\rangle, \qquad |\Psi_N^*\rangle := \sum_{\pi \in \operatorname{LP}_N^*} \Psi_N^*(\pi) |\pi\rangle$$

$$H_N|\Psi\rangle=|\Psi\rangle$$

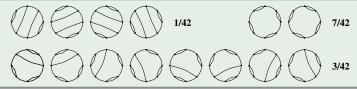




By renormalizing the vectors $|\Psi_n\rangle$ and $|\Psi_N^*\rangle$ [Batchelor, de Gier, Nienhuis, Razumov, Stroganov]

• All the $\Psi_N(\pi)$ are "small" integers.

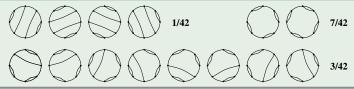
Example (LP₈)



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- Their sum is equal to the enumeration A(n) of Alternating Sign Matrices of size n for $\pi \in LP_{2n}$, or to the enumeration $A_{HT}(N)$ of Half-Turn Symmetric Alternating Sign Matrices of size N for $\pi \in LP_N^*$.

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Does each component have a combinatorial interpretation?

R-matrix

$$\hat{X}_i(z) = \frac{qz - q^{-1}}{q - q^{-1}z}\mathbf{1} + \frac{z - 1}{q - q^{-1}z}\mathbf{e}_i, \qquad \tau = -q - q^{-1}$$

Yang-Baxter equation

$$\hat{X}_{i}(z_{2})\hat{X}_{i+1}(z_{1}z_{2})\hat{X}_{i}(z_{1}) = \hat{X}_{i+1}(z_{1})\hat{X}_{i}(z_{1}z_{2})\hat{X}_{i+1}(z_{2})$$

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Scattering matrices

$$S_i(\vec{z}) = \hat{X}_{i-2}(z_i/z_{i-1})\hat{X}_{i-3}(z_i/z_{i-2})\cdots\hat{X}_{i+1}(z_i/z_{i+2})\hat{X}_i(z_i/z_{i+1})$$

At $q=e^{2\pi i/3}$ (au=1), the scattering equations [Di Francesco, Zinn-Justin]

$$|S_i(\vec{z})|\Psi(\vec{z})\rangle = R^{-1}|\Psi(\vec{z})\rangle$$

have a unique solution (up to normalization), polynomial in \vec{z} .



[P. Di Francesco, P. Zinn-Justin]

At $z_i=1$ the vector $|\Psi(\vec{z})\rangle$ reduces to the stationary probability of the T-L Stochastic model

$$|\Psi_n(\vec{1})\rangle = |\Psi_n\rangle, \qquad |\Psi_N^*(\vec{1})\rangle = |\Psi_N^*\rangle$$

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- Relation with the geometry of Orbital varieties.

Mills, Robbins and Rumsey's extension of Dodgson (aka Lewis Carroll) condensation algorithm ['83]

 $\det\ M\ \det\ M_{1,n}^{1,n}\ =\ \det\ M_n^n\ \det\ M_1^1\ -1\ \det\ M_1^n\ \det\ M_n^1$



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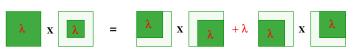






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The result is (surprisingly) a Laurent polynomial in entries m_{ij} :

$$\det_{\lambda} M = \sum_{B \in ASM_n} \lambda^{\prime(B)} (1 + \lambda^{-1})^{N(B)} \prod_{i,j} m_{i,j}^{B_{i,j}}$$

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 $\mathsf{ASM} = \mathsf{Alternating} \ \mathsf{Sign} \ \mathsf{Matrices}$

. . . The Laurent phenomenon is well known in the context of Hirota equations and has led to the development of Fomin-Zelevinsky Cluster Algebras

Alternating Sign Matrices [Mills, Robbins, Rumsey]

Square $n \times n$ matrices with entries 0, 1, -1, such that

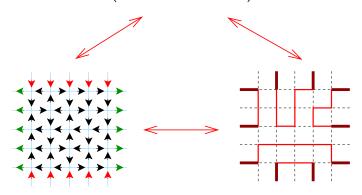
- signs +1 and -1 alternate on each row and each column;
- each row and each column sums to 1.

$$\begin{pmatrix} 0 & 1 & 0 & 0 & 0 \\ 1 & -1 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 \\ 0 & 1 & -1 & 1 & 0 \\ 0 & 0 & 1 & 0 & 0 \end{pmatrix}$$

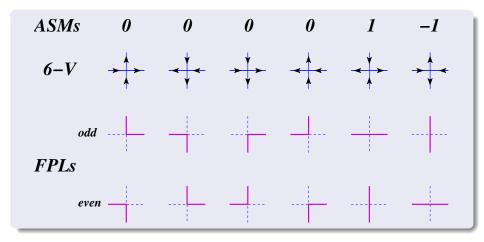
Enumeration
$$A_n = \prod_{j=0, n-1} \frac{(3j+1)!}{(n+j)!}$$
 [Zeilberger '95].

Simpler proof by Kuperberg ['96]: use equivalence with 6-vertex

$$\begin{pmatrix} 0 & 1 & 0 & 0 & 0 \\ 1 & -1 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 \\ 0 & 1 & -1 & 1 & 0 \\ 0 & 0 & 1 & 0 & 0 \end{pmatrix}$$

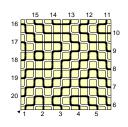


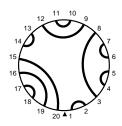
$\overline{\mathsf{ASM}} \leftrightarrow \mathsf{6}\text{-}\mathsf{Vertex} \ (\mathsf{DWBC}) \leftrightarrow \mathsf{FPL}$



Refined FPL enumerations

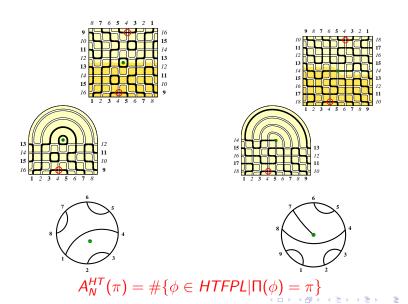
Enumerations of FPLs whose boundary points have a given connection pattern π





$$A_n(\pi) = \#\{\phi \in FPL | \Pi(\phi) = \pi\}$$

Half-Turn ASM



The Razumov Stroganov ex-conjecture

The number $A_n(\pi)$ are related to stationary probability of the Temperley-Lieb stochastic process

Theorem: R-S ex-conjecture '01 [L.C., A. Sportiello '10]

$$\Psi_n(\pi) = \frac{A_n(\pi)}{\sum_{\pi} A_n(\pi)}, \qquad \Psi_N^*(\pi) = \frac{A_N^{HI}(\pi)}{\sum_{\pi} A_N^{HT}(\pi)}.$$

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One can obtain a lot of non trivial informations on the properties of the combinatorial objects on the right-hand side, by studying the left-hand side.

- New proof of the enumeration formula of alternating sign matrices
- Unified proof of enumeration formula of ASM with symmetries
- Number of FPL belonging to certain classes can be counted

Rotational invariance

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Theorem (Wieland, '00)

The enumerations $A_n(\pi)$ and $A_N^{HT}(\pi)$ are invariant under cyclic rotations

$$A_n(\pi) = A_n(R \circ \pi), \qquad A_N^{HT}(\pi) = A_N^{HT}(R \circ \pi)$$
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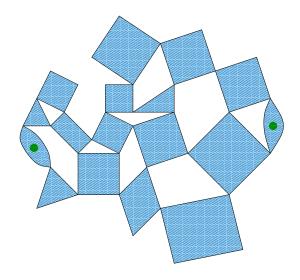
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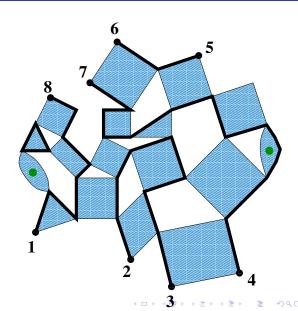
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Let's see how the proof works: this will provide a crucial tool for the proof of the RS conjecture.

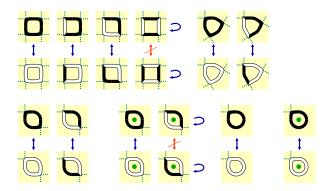
- Consider a planar graph, which is obtained by gluing at the corners 2-, 3- or 4-gons. Inside a 2-gon we can place a puncture.
- FPL = coloring of the edges such that a vertex of cordination 4 is adjacent to 2 colored edges



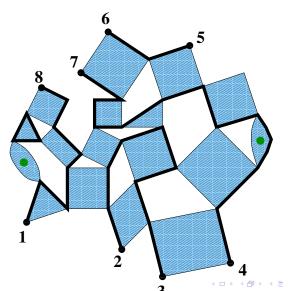
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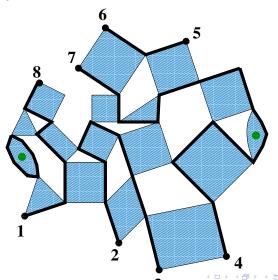
Define the following operation, called Half-Gyration, which preserves the FPL condition at each vertex:



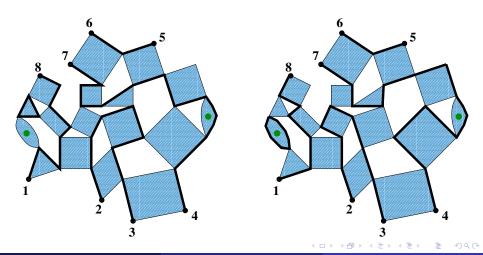
Take an FPL

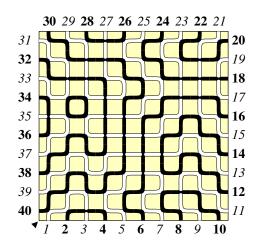


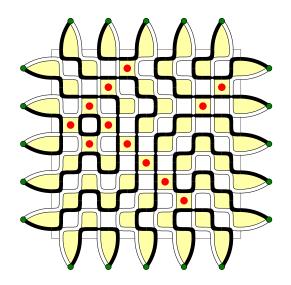
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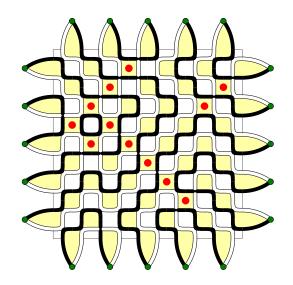


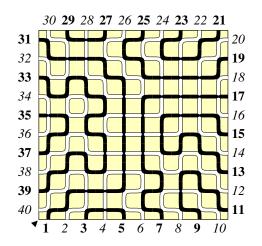
The connectivities of the "boundary" points and the "topological" location of the puncture are preserved

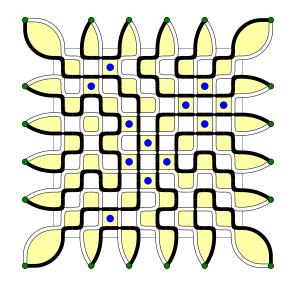


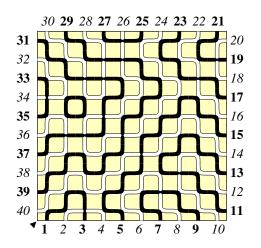


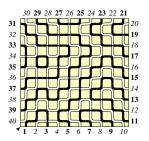


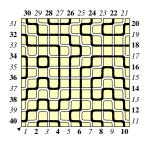












Rotational invariance: Dihedral Domains

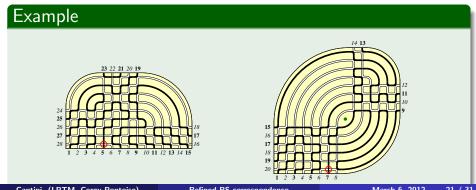
The invariance under rotation of the FPL enumerations is valid on more general planar domains that we call *Dihedral Domains*:

- the bulk of the dual graph must be bipartite,
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Generalized Razumov-Stroganov correspondence

The Razumov-Stroganov correspondence remains valid on these more general domains!

Define again $A_D(\pi)$ as the number of FPL on the domain D, form the vector

$$|\Psi_D^{FPL}
angle = \sum_{\pi} A_D(\pi) |\pi
angle$$

it satisfies [L.C., A. Sportiello, '10]

$$\sum_{i=1}^{2n}(e_i-1)|\Psi_D^{FPL}
angle=0$$

In particular we have proportionality of the enumerations corresponding to the same link pattern on different domains

$$A_D(\pi) = K_D A_n(\pi) \quad \forall \pi$$

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- The proof is valid on all the Dihedral Domains.
- The proportionality factor K_D is equal to the number of FPL corresponding to "Rainbow" Link Patterns: and has often an alternative combinatorial interpretation as the number of dimer covering of regions of the hexagonal lattice.
- Alternative way to compute the total enumerations for FPL on several different classes of domains known from Kuperberg and new ones: for example Quarter Turn Symmatric ASMs of size 4n.

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• Parametrize the Boltzmann weights of the 6-vertex model in terms of spectral parameters: partition function (IK determinant) matches $\sum_{\pi} \Psi_n(\pi, \vec{z})$ but *doesn't work* component-wise!

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- Parametrize the Boltzmann weights of the 6-vertex model in terms of spectral parameters: partition function (IK determinant) matches $\sum_{\pi} \Psi_n(\pi, \vec{z})$ but *doesn't work* component-wise!
- On the other side Di Francesco noticed that $\Psi(\pi;t) := \Psi(\pi,z_1 = \frac{qt+q^{-1}}{q+q^{-1}t},\vec{1})$ are polynomials in t with positive integer coefficients.

$$egin{aligned} \Psi_1(t) &= \{1\} \ \Psi_2(t) &= \{1,t\} \ \Psi_3(t) &= \{1+t,1,t,t(1+t),t^2\} \ \Psi_4(t) &= \{2+3t+2t^2,1+2t,1+t+t^2,2+t,1,t(2+t),t,t(1+2t),t(2+3t+2t^2),t(1+t+t^2),t^2,t^2(2+t),t^2(1+2t),t^3\} \end{aligned}$$

Do these coefficients have a combinatorial meaning?

Di Francesco tried to compare them with enumerations of FPLs weighted by the position $h(\phi)$ of the straight tile on the last row



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The weighted enumerations can be gathered in a vector

$$A_n(\pi;t) := \sum_{\phi \mid \Pi(\phi) = \pi} t^{h(\phi)-1}, \qquad |\Psi_n^{FPL}(t)\rangle := \sum_{\pi} A_n(\pi;t) |\pi\rangle.$$

Unfortunately this doesn't match with $|\Psi_n(t)\rangle$

$$|\Psi_n^{FPL}(t)\rangle \neq |\Psi_n(t)\rangle$$
 !!!!

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Conjecture [Di Francesco '04]

$$\operatorname{Sym}|\Psi_n^{FPL}(t)
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$$\operatorname{Sym}|\Psi_{D}^{FPL}(t)\rangle = K_{D}(t)\operatorname{Sym}|\Psi_{n}(t)\rangle \quad \text{with} \quad \operatorname{Sym} = \sum_{i=1}^{2n} R^{i}$$

but it works also on any Dihedral Domain D!!!.



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Best possible hope:

- Maybe Di Francesco's way to associate a weight or even of associating a *link pattern* to an FPL is only "almost right".
- ullet There is a new way $\tilde{\pi}(\phi)$ of associating link patterns to FPL such that

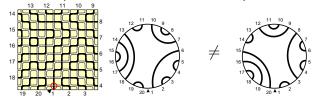
$$| ilde{\Psi}_{D}^{FPL}(t)
angle \propto |\Psi_{n}(t)
angle$$

with no need of symmetrization.

The improved refinement

Here is the rule:

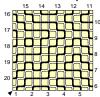
If the refinement position is odd, just start the counting of the external points from the refinement position



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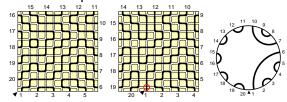
If the refinement position is even: swap the colorations of the edges and then start the counting of the external points from the refinement position:



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The scattering equation

It is not difficult to show that $|\Psi_n(t)\rangle$ is determined (up to normalization) by the scattering equation

$$(\hat{X}_1(t)-R)|\Psi_n(t)
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 $X_1(t)=t+(1-t)e_1$

Theorem [L.C., A. Sportiello '12]

The vector $|\tilde{\Psi}_D^{\mathit{FPL}}(t)
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While this theorem is stronger than the Razumov Stroganov original conjecture, its proof is much simpler!!!!

The really non-trivial "work" was in finding the right way to associate a link pattern to an FPL!

The idea of the proof

Apply the two projectors \mathbf{e}_1 and $(\mathbf{1}-\mathbf{e}_1)$, to the Scattering equation

$$(\mathbf{1} - \mathbf{e}_1)(t\mathbf{1} - R)|\tilde{\Psi}(t)_D^{FPL}
angle = 0, \qquad (\mathbf{e}_1 - R\mathbf{e}_N)|\tilde{\Psi}_D^{FPL}(t)
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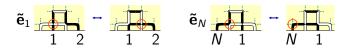
• For the left equation we have just to check that if $\pi \notin \operatorname{Im} \mathbf{e}_1$, then $t\Psi(\pi,t)=\Psi(R^{-1}\pi,t)$. Just a Half-Gyration provides the bijection between FPLs associated to π and $R^{-1}\pi$.

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- ullet For the right equation we notice that, while we don't know how to "act" with all the TL generators on a FPL we know how to act with $ullet_1$ and $ullet_N$



which, combined with two Half-Gyrations, provide the bijection we want.

Back to Di Francesco's conjecture

Corollary: Di Francesco's conjecture [L.C., A. Sportiello '12]

$$\mathrm{Sym}| ilde{\Psi}_D^{FPL}(t)
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Back to Di Francesco's conjecture

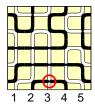
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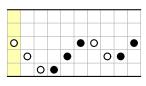
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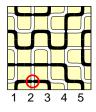
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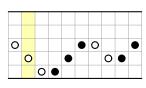




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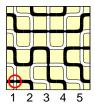
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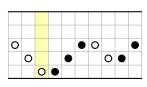




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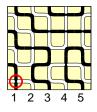
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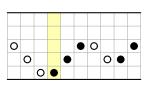




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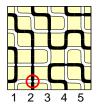
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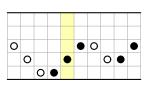




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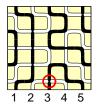
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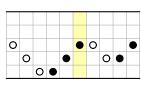




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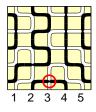
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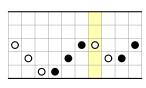




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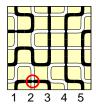
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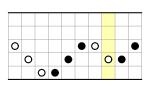




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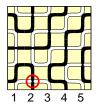
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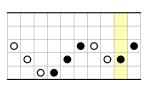




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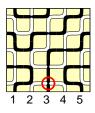
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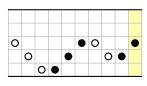




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More ambitious: what about the Razumov Stroganov conjectures without cyclic invariance?

Hint for the proof could be to find the class of domains on which these conjectures hold.