# Lepton Flavour Violating Z decays @ LHC

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- 1. preliminaries: Zs at the LHC
- 2. introduction:
  - what is LFV? and why is it interesting?
  - some upper bounds
- 3. effective operators for  $Z \to \ell \overline{\ell'}$ ...
  - an improbable place to find LFV
- 4.  $Z \to \tau^{\pm} \mu^{\mp}$  @ LHC ... in the  $Z \to \tau^{+} \tau^{-}$  background?
- 5. real work (thanks to S Lacroix)
- 6. expected limits

#### the LHC is not the wrong place to do Z physics

- LEP1 was a clean Z machine, with  $17 \times 10^6 Zs$  $BR(Z \to e^{\pm}\mu^{\mp}) < 1.7 \times 10^{-6}$ ,  $BR(Z \to e^{\pm}\tau^{\mp}) < 9.8 \times 10^{-6}$ ,  $BR(Z \to \mu^{\pm}\tau^{\mp}) < 1.2 \times 10^{-5}$
- at the 7,8 TeV LHC,  $\sigma(pp \rightarrow Z \rightarrow \mu \bar{\mu}) \sim \text{nb.} \ \mathcal{L} \sim 20 \text{ fb}^{-1} \Rightarrow 20 \times 10^6 \text{ Zs ??}$

$$\#Zs \simeq \frac{\sigma(pp \to Z \to \mu\bar{\mu}) \times \mathcal{L}}{BR(Z \to \mu\bar{\mu})} \sim 10^8 Zs$$

 $BR(Z \to \mu \bar{\mu}) \simeq 0.0366$ 

(compare e.g.  $\sigma(pp \to t\bar{t}) \sim 160 \text{ pb...} \gtrsim 150 \text{ Zs}$  for each  $t\bar{t}$  pair)

#### Lepton Flavour Violation: what is it? Why interesting?

 ${\sf LFV} \equiv {\sf flavour\ changing\ point\ interaction\ of\ charged\ leptons} \\ \equiv {\sf FCNC\ in\ charged\ leptons}:\ \tau \to \mu\gamma,\ \dots$ 

- 1. we know  $m_{\nu} \neq 0 \Rightarrow Beyond the Standard Model in the leptons!$
- 2. But not see LFV yet.



3. But  $A(LFV) \propto m_{\nu}^2/m_W^2 \sim 20^{-24}$ ,  $\Rightarrow$  observable LFV requires dynamics other than  $m_{\nu}$ 

entertainment for theorists: obtain log GIM in leptons...

### What do we know (experimentally)

some processes	current sensitivities
$BR(\mu \to e\gamma)$	$< 2.4 \times 10^{-12}$
$BR(\mu \to e\bar{e}e)$	$< 1.0 \times 10^{-12}$
$\frac{\sigma(\mu + Au \rightarrow e + Au)}{\sigma(\mu \text{ capture})}$	$< 7 \times 10^{-13}$
$BR( au  o \ell \gamma)$	$< 3.3, 4.4 \times 10^{-8}$
$BR( au  ightarrow 3\ell)$	$< 1.5 - 2.7 \times 10^{-8}$
$BR(\tau \to e\phi)$	$< 3.1 \times 10^{-8}$
$BR(\overline{K_L^0} \to \mu \bar{e}) \\ BR(K^+ \to \pi^+ \bar{\nu} \nu)$	$< 4.7 \times 10^{-12}$ = 1.7 ± 1.1 × 10 <sup>-10</sup>
$BR(Z \to e^{\pm} \mu^{\mp})$	$1.7 \times 10^{-6}$
$BR(Z \to e^{\pm}\tau^{\mp})$ $BR(Z \to \mu^{\pm}\tau^{\mp})$	$9.8 \times 10^{-6}$ $1.2 \times 10^{-5}$

#### How to interpret those numbers —two perspectives

1.  $m_{\nu}$  arise in my favourite model — what do LFV bounds tell me about it?

2. I want to know what generates  $m_{\nu}$  — how do I learn that from the data?

? maybe I learn something with the Effective Lagrangian?

#### (Organising and interpreting) what we know: the effective Lagrangian

Suppose that NP<sub>(articles)</sub> are above (fuzzy) mass scale  $M > m_Z$ . At  $E \ll M$ , describe their effects as "contact interactions" among light particles in an "effective Lagrangian":

$$\mathcal{L}_{eff} = \mathcal{L}_{SM} + \Delta \mathcal{L}_{eff}^{LFV} + \Delta \mathcal{L}_{eff}^{other}$$
$$\Delta \mathcal{L}_{eff}^{LFV} = \sum_{d \ge 5} \sum_{n} \frac{C^n}{M^{d-4}} O_n(H, \{\psi\}, A_\mu, ...) + h.c.$$

The operators  $\{O_n\}$  :

- built with \*kinematically accessible\* SM fields (avec Z, at  $m_Z$ ; sans Z à  $m_\tau$ )
- respect SM gauge symmetries
- describe the legs of the LFV diagrams (including Higgs vevs)

The (dimless) coefficients  $C^n$  contain coupling constants,  $1/16\pi^2$ , ...

(SM = Standard Model, NP = New Physics, New Particles +...)

#### At dimension six in $\mathcal{L}_{eff}$ , at scales $\lesssim m_Z$

match in two steps: 1) @ M: EW + NP onto broken SM with particles  $\{Z, W^{\pm}, \tau, \mu, e, \nu_{\alpha}, \gamma\}$  and  $\mathcal{L}_{eff}$ 2) @  $m_Z(m_W)$ : onto SM with particles  $\{\tau, \mu, e, \nu_{\alpha}, \gamma\}$  and  $\mathcal{L}'_{eff}$ After step 1):



(NB I assume NP in loops;  $\propto 1/(16\pi^2 M^2)$ , for most LFV processes)

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For a given process with  $BR < \dots$ , can obtain a lower bound on M:

- 1. identify operators/diagrams corresponding to a process,
- 2. set  $C \simeq 1$ ,
- 3. compute rate,...

### Interpreting what we know: bounds assuming dimension 6 operators

process	bound	scale, dim 6, loop
$BR(\mu \to e\gamma)$ $BR(\mu \to e\bar{e}e)$ $\frac{\sigma(\mu + Ti \to e + Ti)}{\sigma(\mu Ti \to \nu Ti')}$	$< 2.4 \times 10^{-12}$ $< 1.0 \times 10^{-12}$ $< 4.3 \times 10^{-13}$	
$BR(\tau \to \ell \gamma) BR(\tau \to 3\ell) BR(\tau \to e\pi)$	$< 3.3, 4.4 \times 10^{-8}$ $< 1.5 - 2.7 \times 10^{-8}$ $< 8.1 \times 10^{-8}$	
$BR(\overline{K^0_L} \to \mu \bar{e})$	$< 4.7 \times 10^{-12}$	
$\begin{array}{l} BR(Z \rightarrow e^{\pm} \mu^{\mp}) \\ BR(Z \rightarrow e^{\pm} \tau^{\mp}) \\ BR(Z \rightarrow \mu^{\pm} \tau^{\mp}) \end{array}$	$< 1.7 \times 10^{-6}$ $< 9.8 \times 10^{-6}$ $< 1.2 \times 10^{-5}$	0.22 TeV 0.14 TeV 0.14 TeV

can produce such NP at LHC? EFT marginally consistent?

#### At dimension six in $\mathcal{L}_{eff}$ , at scales $\lesssim 10$ GeV



For a given process with  $BR < \dots$ , can obtain a lower bound on M:

1. identify operators/diagrams corresponding to a process,

2. set  $C \simeq 1$ ,

3. compute rate,...

#### Interpreting what we know: bounds assuming dimension 6 operators

process	bound	scale, dim 6, loop
$BR(\mu \to e\gamma)$	$< 2.4 \times 10^{-12}$	$48 { m TeV}$
$BR(\mu \to e\bar{e}e)$	$< 1.0 \times 10^{-12}$	174  TeV (tree)
		$14 { m TeV}$
$\frac{\sigma(\mu + Ti \rightarrow e + Ti)}{\sigma(\mu Ti \rightarrow \nu Ti')}$	$< 4.3 \times 10^{-13}$	40 TeV
$BR(\tau \to \ell \gamma)$	$  < 3.3, 4.4 \times 10^{-8}$	2.8 TeV
$BR( au  ightarrow 3\ell)$	< 1.5 - 2.7 × 10 <sup>-8</sup>	$0.8 { m TeV}$
$BR(\tau \to e\pi)$	$< 8.1 \times 10^{-8}$	0.5 TeV
$BR(\overline{K_L^0} \to \mu \bar{e})$	$< 4.7 \times 10^{-12}$	$25 \text{ TeV}(V \pm A)$
		140 TeV $(S \pm P)$
$BB(7 \rightarrow e^{\pm}u^{\pm})$	$< 1.7 \times 10^{-6}$	0.22  To/
$DR(Z \to e^+ \mu^+)$ $DR(Z \to e^+ e^{\pm})$	$  1.1 \times 10  $	0.22  TeV
$BK(Z \to e^{\pm}\tau^{+})$	$< 9.8 \times 10^{\circ}$	0.14 IeV
$BR(Z \to \mu^{\perp} \tau^{+})$	$< 1.2 \times 10^{-5}$	0.14 leV

if all flavour-changing couplings are of the same order, then should look for LFV in  $\mu \rightarrow e$ 

lepton decays probe higher M that Z decays— in EFT, given  $\mu, \tau$  bounds, can LFV Z decay be observed?

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Mass dimension of Z and two lepton external legs = 4  $\Rightarrow Z \rightarrow \tau^{\pm} \mu^{\mp}$  operators contains two Higgs and/or Derivatives Three options among gauge invariant operators at dimension 6:

 $\mathcal{O}(\partial^2) : \overline{\mu}\gamma_{\beta}D_{\alpha}\tau B^{\alpha\beta} , \dots$  $\mathcal{O}(H^2) : [H^{\dagger}D_{\alpha}H]\overline{\mu}\gamma^{\alpha}\tau , \dots$  $\mathcal{O}(yH\partial) \text{ dipole } : \overline{\ell}_{\mu}H\sigma_{\beta\alpha}\tau B^{\alpha\beta} , \dots$ 

(where  $B^{\alpha\beta} = \partial^{\alpha}B^{\beta} - \partial^{\beta}B^{\alpha}$ , B hypercharge gauge boson).

Mass dimension of Z and two lepton external legs = 4 $\Rightarrow$  operator contains two Higgs and/or Derivatives Three options among gauge invariant operators at dimension 6:

 $\mathcal{O}(\partial^2) \quad : \quad \overline{\mu}\gamma_\beta D_\alpha \tau B^{\alpha\beta} \quad , \quad \overline{\ell}_\mu \sigma^I \gamma_\beta D_\alpha \ell_\tau W^{I\alpha\beta} \quad , \quad \overline{\ell}_\mu \gamma_\beta D_\alpha \ell_\tau B^{\alpha\beta}$ 

 $\mathcal{O}(H^2) : [H^{\dagger}D_{\alpha}H]\overline{\mu}\gamma^{\alpha}\tau , [H^{\dagger}\sigma^{I}D_{\alpha}H][\overline{\ell}_{\mu}\sigma^{I}\gamma^{\alpha}\ell_{\tau}] , [H^{\dagger}D_{\alpha}H][\overline{\ell}_{\mu}\gamma^{\alpha}\ell_{\tau}]$  $\mathcal{O}(yH\partial) \text{ dipole } : \overline{\ell}_{\mu}H\sigma_{\beta\alpha}\tau B^{\alpha\beta} , \overline{\ell}_{\mu}\sigma^{I}H\sigma_{\beta\alpha}\tau W^{I\alpha\beta}$ 

Need two powers of a vev/momentum in operator. Three options among gauge invariant operators at dimension 6. Suppose operator coefficients such that:

Rossi+Brignole

$$\dots, \ \overline{\mu}\gamma_{\beta}D_{\alpha}\tau B^{\alpha\beta} \rightarrow g_{Z}C\frac{p_{Z}^{2}}{16\pi^{2}M^{2}}\overline{\mu}\gamma_{\alpha}\tau Z^{\alpha}$$
$$\dots, \ [H^{\dagger}D_{\alpha}H]\overline{\mu}\gamma^{\alpha}\tau \rightarrow g_{Z}A\frac{m_{Z}^{2}}{16\pi^{2}M^{2}}\overline{\mu}\gamma_{\alpha}Z^{\alpha}\tau$$
$$\dots, \ \overline{\ell}_{\mu}H\sigma_{\beta\alpha}\tau B^{\alpha\beta} \rightarrow g_{Z}D\frac{m_{\tau}}{16\pi^{2}M^{2}}[\overline{\mu}\sigma_{\alpha\beta}\tau]Z^{\alpha\beta}$$

NP of mass  $M > m_Z$  in a loop, A, C, D dimless

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NP of mass  $M > m_Z$  in a loop, A, C, D dimless

Operators with gradients better constrained at higher energies:

on the 
$$Z$$
: vertex  $=g_Z \frac{Cm_Z^2}{16\pi^2 M^2} \overline{\mu} Z \tau$ ,  $BR(Z \to \tau^{\pm} \mu^{\mp}) \sim 1.7 \times 10^{-5} \frac{m_Z^4}{M^4}$ ,  $(C = 1)$ 

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NP of mass  $M > m_Z$  in a loop,  $A, C, D \mbox{ dimless}$ 

Operators with gradients better constrained at higher energies:

on the 
$$Z$$
: vertex  $=g_Z \frac{Cm_Z^2}{16\pi^2 M^2} \overline{\mu} \mathbb{Z} \tau$ ,  $BR(Z \to \tau^{\pm} \mu^{\mp}) \sim 1.7 \times 10^{-5} \frac{m_Z^4}{M^4}$ ,  $(C = 1)$   
in  $\tau^{\pm} \to \mu^{\mp} \mu^{\pm} \mu^{\pm}$ : vertex  $< g_Z \frac{Cm_{\tau}^2}{16\pi^2 M^2} \overline{\mu} \mathbb{Z} \tau$   
 $\mathcal{T}_L$   
 $\mathcal{T$ 

### The gradient<sup>2</sup> $Z \rightarrow \tau^{\pm} \mu^{\mp}$ operators: are they important in loops?

and can I calculate that?

 $\begin{array}{c} Z \\ \tau \\ \tau \\ \gamma \end{array}$ 

- 1. assume NP scale  $M \gg m_Z$
- 2. assume NP generates only  $\partial^2$  operator (no other LFV; not  $\tau \to \mu\gamma$ ), so "interaction":

$$g_Z C_{\mu\tau} \frac{p_Z^2}{16\pi^2 M^2} \overline{\mu} \gamma_\alpha \tau Z^\alpha$$

3. in RG running between M and  $m_Z$ ,  $Z \to \tau^{\pm} \mu^{\mp}$  will mix to  $\tau \to \mu \gamma$  operator (...estimate the coefficient of  $1/\epsilon$  in dim reg...)

$$\widetilde{BR}(\tau \to \mu \gamma) \simeq \frac{3\alpha}{4\pi} \frac{g_Z^4}{G_F^2 M^4} \left(\frac{C_{\mu\tau} \log}{32\pi^2}\right)^2 \sim 4 \times 10^{-8} \frac{C_{\mu\tau}^2 v^4}{M^4}$$

$$\Rightarrow \text{ no constraint from } \tau \to \ell \gamma$$

but  $\mu \to e\gamma$  constrains  $C_{e\mu}$ :  $BR(Z \to e^{\pm}\mu^{\mp}) \lesssim 10^{-10}$ .

### (parenthèse: $H^2$ and dipole operators can be neglected for $Z \to \tau^{\pm} \mu^{\mp}$ )

neglect  $\mathcal{O}(H^2)$  and  $\mathcal{O}(yH\partial)$  operators, because more strictly constrained elsewhere:

• for 
$$[H^{\dagger}D_{\alpha}H]\overline{\mu}\gamma^{\alpha}\tau \to g_Z A \frac{m_Z^2}{16\pi^2 M^2}\overline{\mu}\gamma_{\alpha}Z^{\alpha}\tau$$

$$\frac{BR(Z \to \tau^{\pm} \mu^{+})}{BR(Z \to \mu^{+} \mu^{-})} \simeq \frac{m_{Z}^{4}}{s_{W}^{4}} |A|^{2} \lesssim 1 \quad , \quad \frac{BR(\tau \to 3\mu)}{BR(\tau \to \mu\nu\bar{\nu})} = \frac{m_{Z}^{4}}{M^{4}} |A|^{2} \lesssim 10^{-2}$$

• for  $\overline{\ell}_{\mu}H\sigma_{\beta\alpha}\tau B^{\alpha\beta} \to g_Z D \frac{m_{\tau}}{16\pi^2 M^2} [\overline{\mu}\sigma_{\alpha\beta}\tau] Z^{\alpha\beta}$ 

probably (?), SM gauge invariant operators contribute also to photon dipole...not pay  $m_{\tau}$  factor in  $\tau \to \mu \gamma$ , so better bounds there.

 $\Rightarrow$  better bounds on coefficients of  $H^2$  and dipole operators from lepton precision than Z decay.

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### $Z \rightarrow \tau^{\pm} \mu^{\mp}$ — how to find at the LHC?

• LEP1 bounds, with  $17 \times 10^6 Z$ s  $BR(Z \to e^{\pm}\mu^{\mp}) < 1.7 \times 10^{-6}$ ,  $BR(Z \to e^{\pm}\tau^{\mp}) < 9.8 \times 10^{-6}$ ,  $BR(Z \to \mu^{\pm}\tau^{\mp}) < 1.2 \times 10^{-5}$ 

• 
$$\mu \to e\gamma \Rightarrow BR(Z \to e^{\pm}\mu^{\mp}) \lesssim 10^{-10}$$

- end 2012, a few  $\times 10^8$  Zs at the LHC
- reconstruct  $\mu, e$  with  $.5 \rightarrow$  few % accuracy (...hadronic  $\tau$ s are "difficult"...)

$$\Rightarrow$$
 study  $Z \rightarrow \tau^{\pm} \mu^{\mp} \rightarrow (e^{\pm} \nu \bar{\nu}) \mu^{\mp}$ 

( recall  $BR(\tau \to \ell \nu \bar{\nu}) \simeq 0.176$ )

- can extrapolate to  $Z \to \tau^{\pm} e^{\mp}$ , because soft  $\mu$  easier to find that soft e (see next page)
- ?? how to find in  $Z \to \tau^{\pm} \tau^{\mp} \to (e^{\pm} \nu \bar{\nu}) (\mu^{\mp} \nu \bar{\nu})$  ??  $(BR(Z \to e\mu + 4\nu) \sim 10^{-3})$

$$p_T$$
 of  $e$  and  $\mu$ , for  $Z \to \tau^{\pm} \mu^{\mp} \to (e^{\pm} \nu \bar{\nu}) \mu^{\mp}$ 



#### **Kinematics**

- Want to distinguish  $pp \to Z \to \tau^+ \mu^- \to (e^+ \nu \bar{\nu}) \mu^$ from background  $pp \to Z \to \tau^+ \tau^- \to (e^+ \nu \bar{\nu}) (\mu^- \nu \bar{\nu}) \dots$
- variable that differs for  $p_T$  with e (signal) from  $p_T$  with e and  $\mu$ :

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- variable that differs for  $p_T$  with e (signal) from  $p_T$  with e and  $\mu$ :
- 1. collinear approx for  $\tau$  decay products  $\tau$  boosted :  $\gamma \sim m_Z/(2m_\tau)$ ,  $\Rightarrow$  all  $\tau$  daughters aligned on  $\tau$ :

$$\begin{split} p_{\tau^+} &= p_{e^+} + p_{\nu} + p_{\bar{\nu}} \equiv \alpha p_{e^+} \\ \text{for backgrd, } Z \to \tau^+ \tau^- \to (e^+ \nu \bar{\nu}) (\mu^- \nu \bar{\nu}) \text{ also } p_{\tau^-} = \beta p_{\mu^-} \\ \text{signal} : p_Z^2 - m_\tau^2 &= 2\alpha p_{e^+} \cdot p_{\mu^-} \quad , \quad \text{background} : m_Z^2 - 2m_\tau^2 = 2\alpha \beta p_{e^+} \cdot p_{\mu^-} \end{split}$$

2. Neglect  $p_T$  of Z: signal :  $\alpha |p_{T,e^+}| = |p_{T,\mu^-}|$ , background :  $\alpha |p_{T,e^+}| = \beta |p_{T,\mu^-}|$ 

 $\Rightarrow$  two determinations of  $p_T = p_{T,\nu} + p_{T,\bar{\nu}}$  ( $\alpha$ ), assuming its aligned on  $p_{T,e}$ . Difference is 0 for signal...and not for background.

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#### **Backgrounds and Simulations...**

Some processes which give  $\mu^\pm e^\mp + \ldots$ , and expected number of events for 20 fb $^{-1}$  data

1.  $Z/\gamma^* \to \tau^{\pm}\tau^{\mp} \to \ell^{\pm}\nu\bar{\nu}\ell'^{\mp}\nu\bar{\nu}$  with  $\ell, \ell' = e, \mu$  ( $M_{Z/\gamma^*} > 20$  GeV) 4 800 000

2.  $t\bar{t} \rightarrow b\ell^+ \nu \bar{b}\ell'^- \bar{\nu}$  with  $\ell, \ell' = e, \mu, \tau$  480 000

3. 
$$Wt \rightarrow \ell^{\pm} \nu b \ell'^{\mp} \nu$$
 with  $\ell, \ell' = e, \mu, \tau$  47 000

4. 
$$W^+W^- \to \ell^+ \nu \ell'^- \bar{\nu}$$
 with  $\ell, \ell' = e, \mu, \tau$  120 000

5.  $Z/\gamma^* Z/\gamma^* \rightarrow f\bar{f}f'\bar{f}'$  160 000

(N)NLO/(N)NLL cross-sections from various codes. LO simulation with Pythia. Fast CMS simulation of Delphes (anti- $k_t$  jets of FastJet).

Simulate  $\sim 10 \times$  number of events expected by end 2012 (grid). And  $10^5$  signal evts.

## Looking for $Z \to \tau^{\pm} \mu^{\mp} \dots$

Selection criteria	$N_{backgrd.}$	Signal efficiency (%)
muon, $p_T > 30 {\rm ~GeV}$	43,500	9.4
e, $p_T > 10  { m GeV}$		
OS	42,652	9.4

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no jet with $p_T > 30 \text{ GeV}$	11,358	7.8
$\Delta \phi(e,\mu) > 2.7$	6,850	6.9
$\Delta \phi(e, E_T) < 0.7$	3,763	6.2
$38  GeV < M_{e\mu} < 92   \mathrm{GeV}$	3,201	6.1

Originally 5.5 M SM background events, are left 3201. Of which, 95% are  $Z/\gamma^* \rightarrow \tau^{\pm}\tau^{\mp} \rightarrow \mu^{\pm}e^{\mp}\nu\bar{\nu}$  (see next page).

signal efficiency : 6.1 %. LEP limit  $(BR(Z \rightarrow \tau^{\pm} \mu^{\pm}) < 1.2 \times 10^{-5}) = 489$  signal events.

#### So where are we now? ...



 $\Delta \alpha \to 0$  for neutrino 4-p aligned on e(dashed line is  $Z \to \tau^{\pm} \mu^{\mp}$  at the LEP1 limit: 489 evts. Backgrd = 3201)

### Getting a bound on $BR(Z \rightarrow \tau^{\pm} \mu^{\mp})$ from that plot... statistics

Want to quantify that the simulated background does not look like the signal (significance test)

Have expected background, and signal efficiency.

Assume 3% systematic uncertainty (!)

Compute 95% CL expected limit...using  $CL_s$  ( $\simeq$  value of BR such that should see more events in 95% of cases):

# $BR(Z ightarrow au^{\pm}\mu^{\mp}) < 3.5 imes10^{-6}$

(4 times better than LEP1)

If look for  $BR(Z \to e^{\pm} \mu^{\mp})$  too...



 $BR(Z \to e^\pm \mu^\mp) < 4.1 \times 10^{-7}$ 

### Systematics...



#### Summary

- Neutrinos have mass  $\Leftrightarrow$  there is New Physics dedicated to Lepton Flavour!
- But, no flavour-changing processes observed among charged leptons (yet).
   ⇒ look everywhere!

• @LHC

- can look for New (s)Particles with LFV decays
- can look for LFV with external legs that exist:

$$Z \to \tau^{\pm} \mu^{\mp}, \tau^{\pm} e^{\mp}, \mu^{\pm} e^{\mp}$$

 with data up to 2013, and aggressive systematic error improvement (→ 3%), can improve LEP bounds by factor ~ 4:

 $BR(Z o au^{\pm} \mu^{\mp}) < 3.5 imes 10^{-6} \ , \ BR(Z o e^{\pm} \mu^{\mp}) < 4.1 imes 10^{-7}$ 

(expect sensitivity  $BR(Z \to \tau^{\pm} e^{\mp})$  similar/better than  $Z \to \tau^{\pm} \mu^{\mp}$ . And  $BR(Z \to e^{\pm} \mu^{\mp}) < 10^{-10}$  from  $\mu \to e\gamma$ ).